

Thermal QCD phases by TWEXT collaboration

A. Yu. Kotov



[AYuK, M.P. Lombardo, A. Trunin, Phys.Lett.B 823, 2021]

[AYuK, M.P. Lombardo, A. Trunin, Symmetry 13, 2021]

[AYuK, M.P. Lombardo, A. Trunin, PoS Lattice 2021]

[AYuK, M.P. Lombardo, A. Trunin, in progress]

New trends in Thermal Phases of QCD, Prague, 2023

TWEXT (Twisted Wilson @ EXTreme) Collaboration

- Andrey Yu. Kotov
- Maria Paola Lombardo
- Anton Trunin

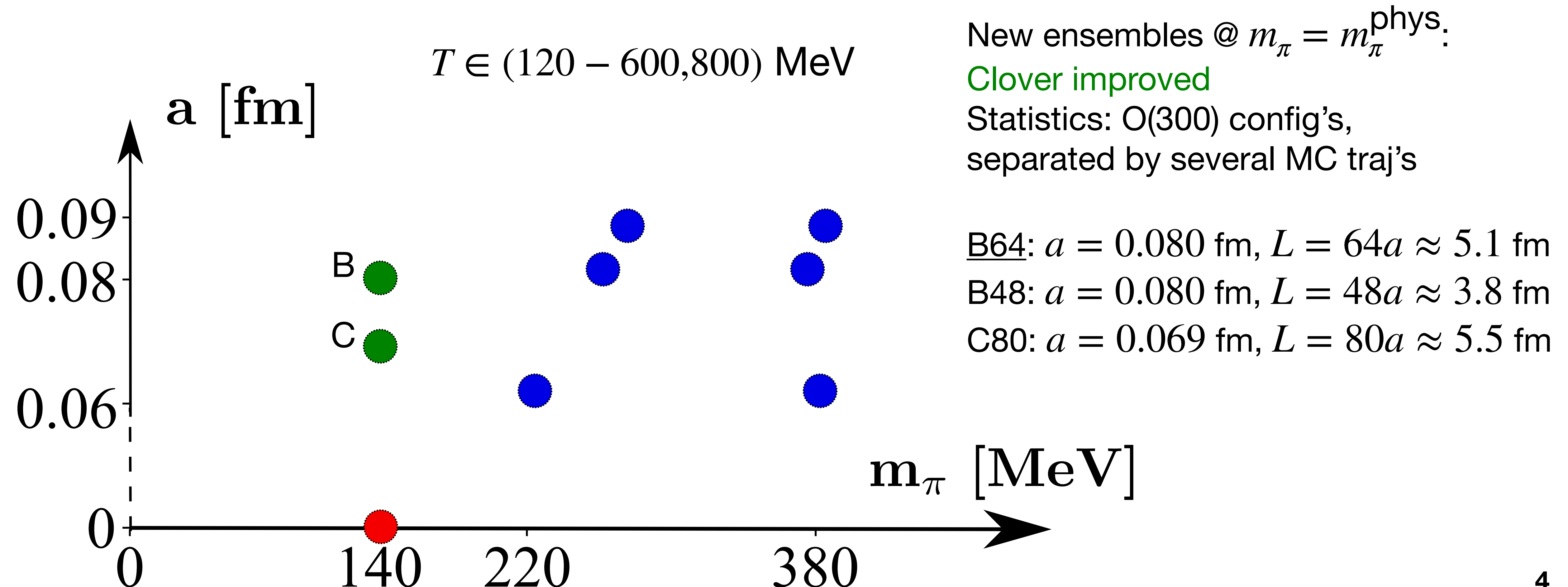
TWEXT (Twisted Wilson @ EXTreme) Collaboration

Lattice details

- 2+1+1 **Wilson twisted mass** fermions at maximal twist
automatically $O(a)$ improved
[R. Frezzotti, G. Rossi, 2004]
- Iwasaki gauge action
- **Heavy quarks** (c, s): close to the physical values
- $m_\pi \in [135, 370]$ MeV
- **Fixed scale approach**: $a = \text{fixed}$, $T \leftrightarrow N_t$ (even for technical reasons)
- Based on ETMC $T = 0$ parameters & tmLQCD code
[C. Alexandrou et al.,2018][C. Alexandrou et al.,2021]

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Ensemble summary



Chiral phase transition & novel order parameter

Novel order parameter

Novel order parameter

- Chiral condensate $\langle \bar{\psi}\psi \rangle$, to be renormalised ($1/a^2$ divergences)

Novel order parameter

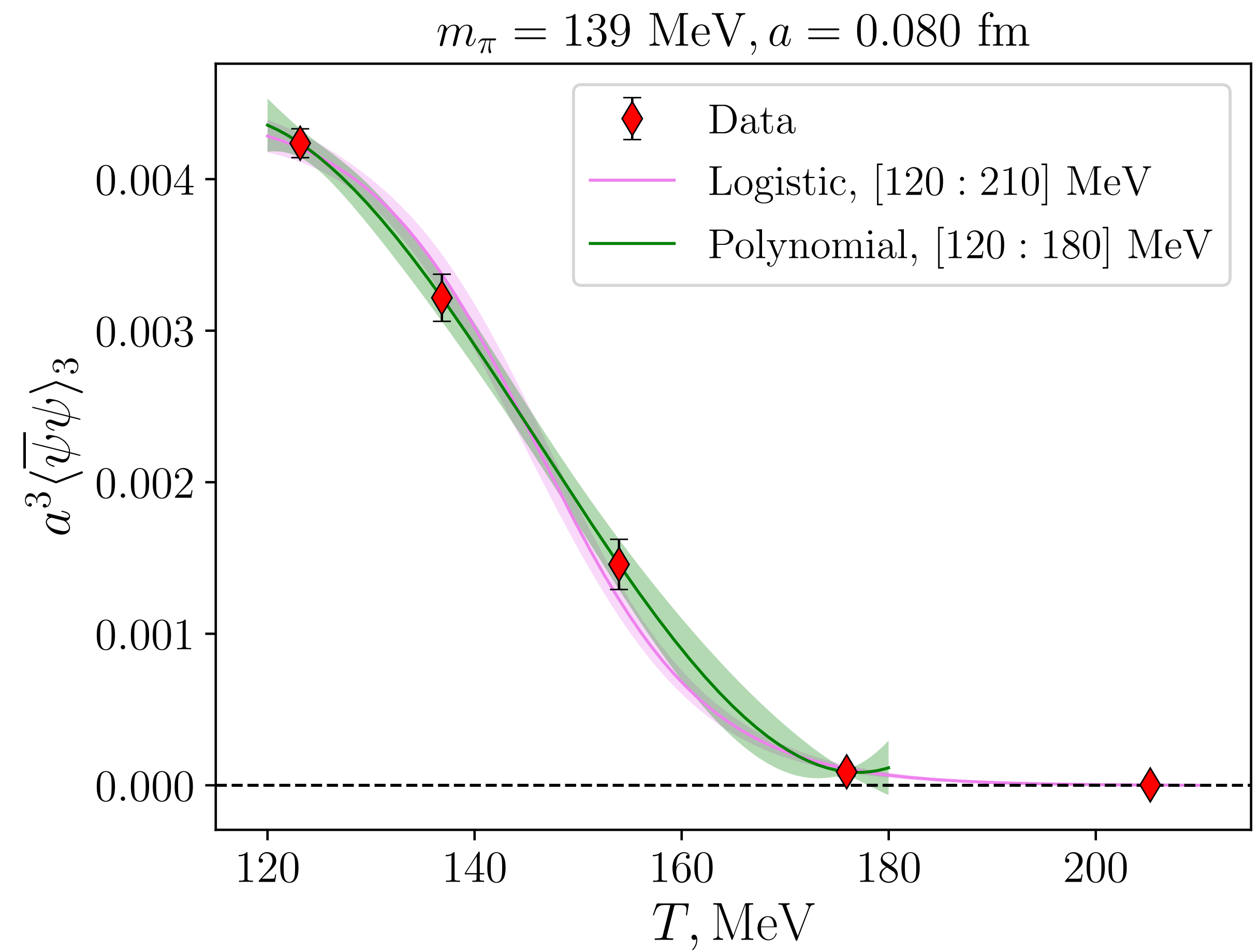
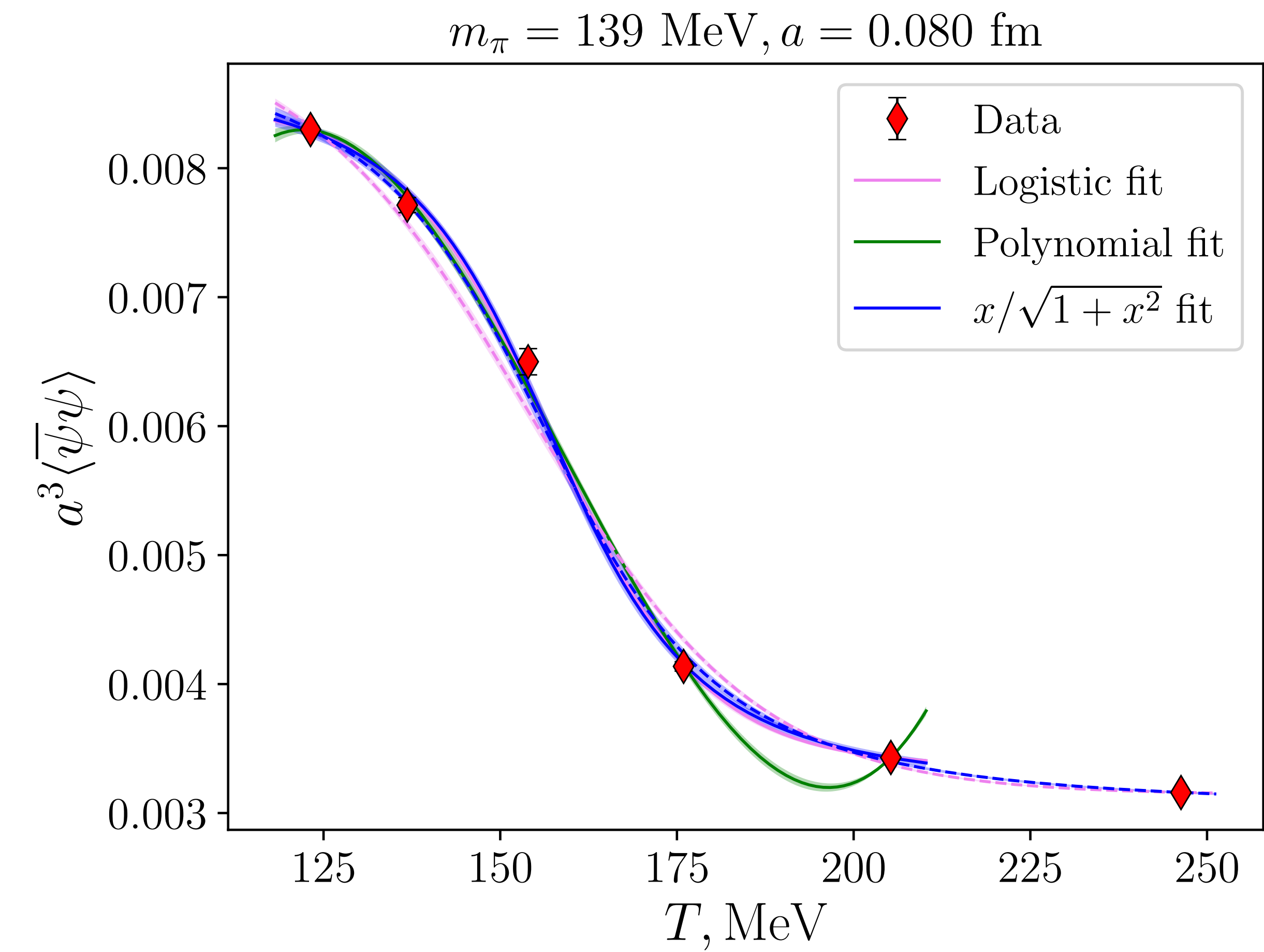
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 - Although not important in the **fixed scale approach**

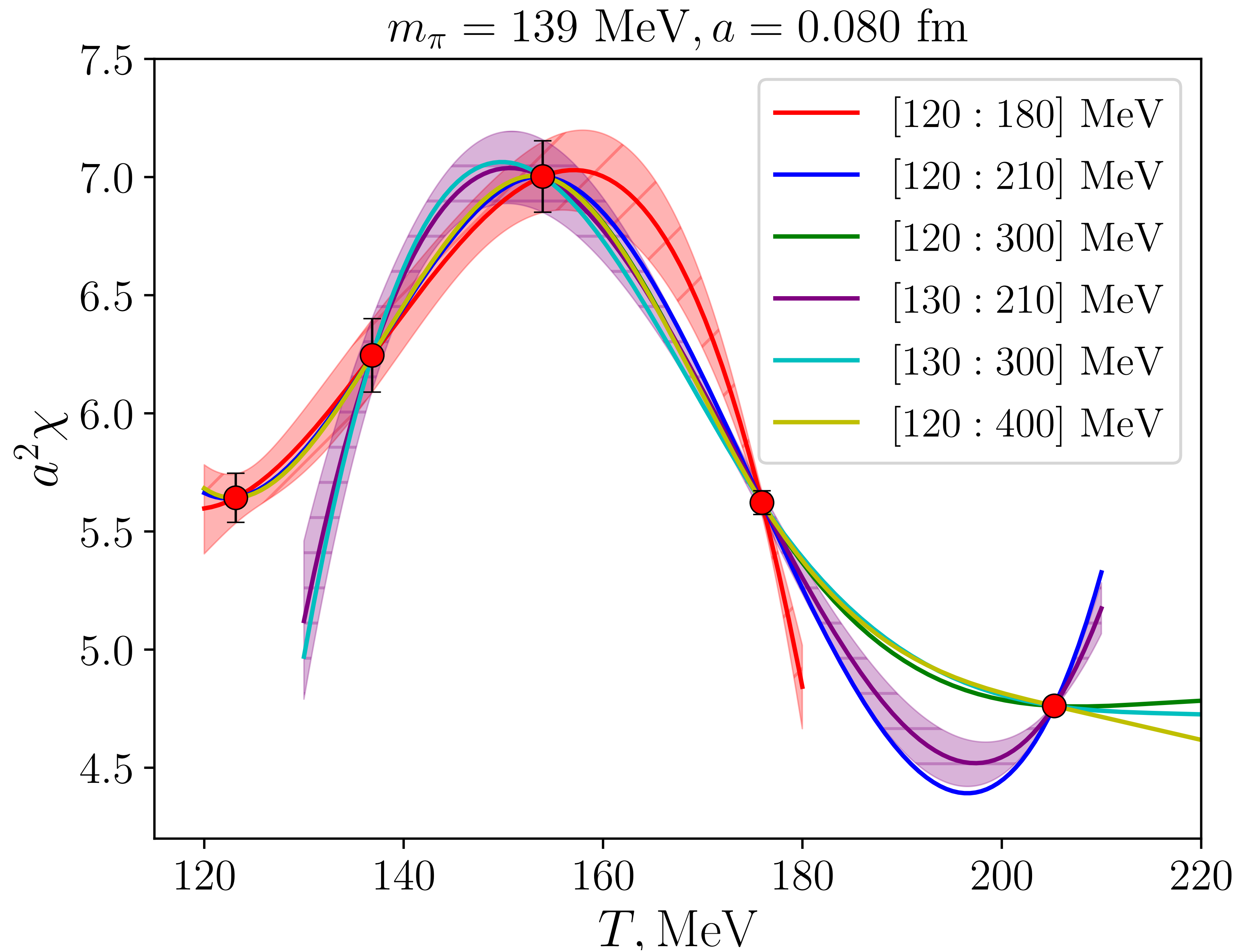
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- Chiral susceptibility $\chi = \partial \langle \bar{\psi}\psi \rangle / \partial m$

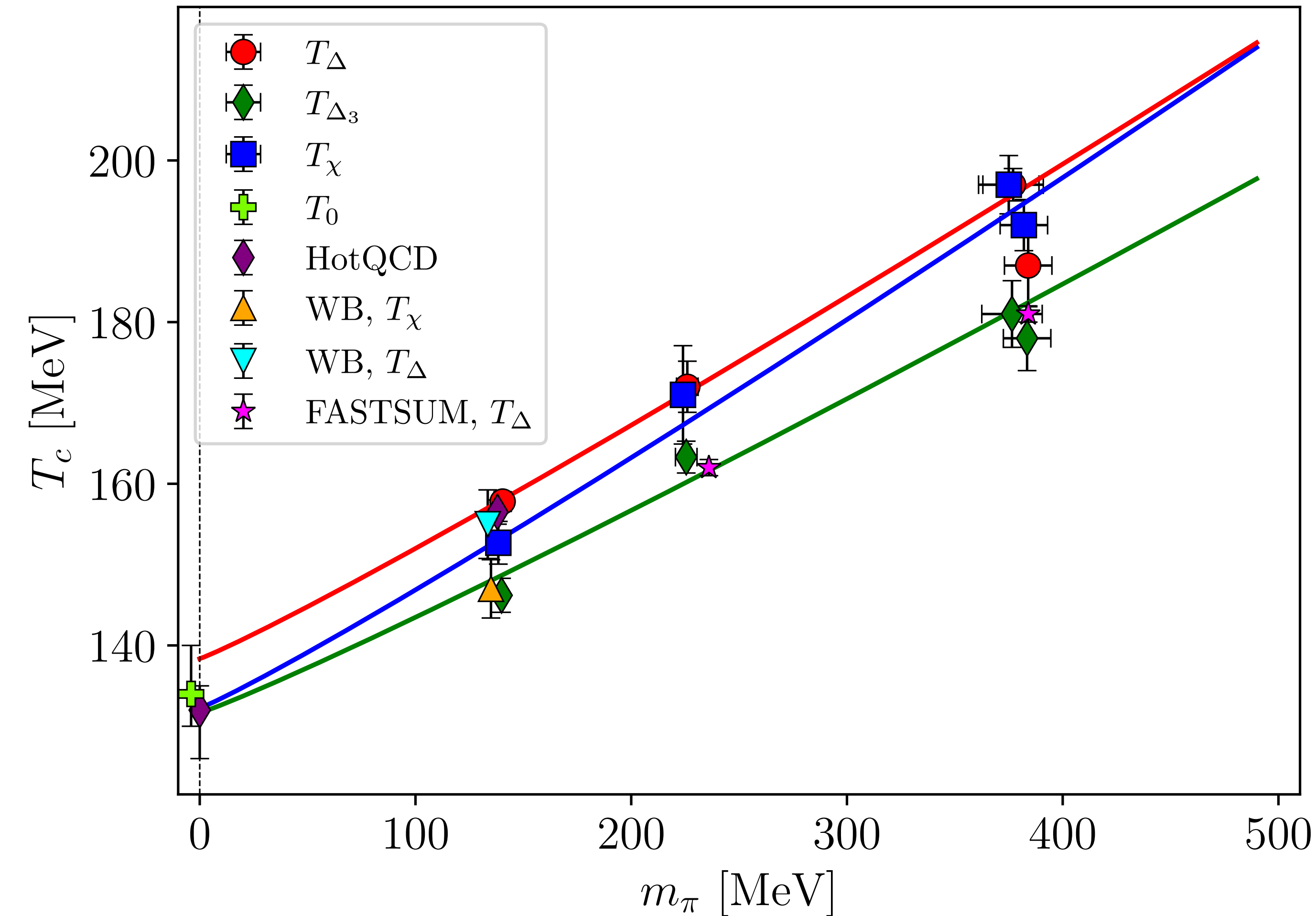
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 - Although not important in the **fixed scale approach**
- Chiral susceptibility $\chi = \partial \langle \bar{\psi}\psi \rangle / \partial m$
- **Novel order parameter:** $\langle \bar{\psi}\psi \rangle_3 = \langle \bar{\psi}\psi \rangle - m\chi$
 - $1/a^2$ divergences cancel
 - $\sim m^3$ (symmetric phase, large T)





Critical temperature and the chiral limit



	$T(m_\pi = 139 \text{ MeV})$ [MeV]	$T(m_\pi = 0)$ [MeV]
$\langle \bar{\psi} \psi \rangle$	157.8(12)	138(2)
χ	153(3)	132(4)
$\langle \bar{\psi} \psi \rangle_3$	146(2)	132(3)

$$T_c = T_c(0) + k_s m_\pi^{2/\beta\delta}, O(4)$$

$$T_0 = 134^{+6}_{-4} \text{ MeV}$$

[AYuK, M.P. Lombardo, A. Trunin, 2021]

Scaling behaviour

Symmetries of QCD with n quarks

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Global symmetry: $U_L(n) \times U_R(n) \cong \boxed{SU_L(n) \times SU_R(n)} \times U_B(1) \times U_A(1)$

Spontaneously broken $\downarrow$ Baryon number Anomalous Broken

$SU_V(n)$

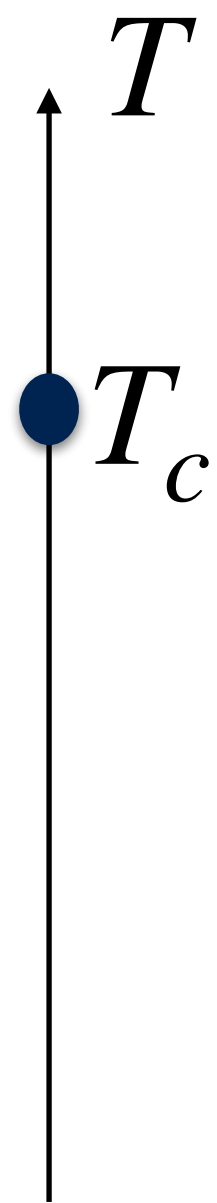
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$T > T_c$ ($m = 0$): (which?) symmetry restoration $\Leftrightarrow$ order (universality)

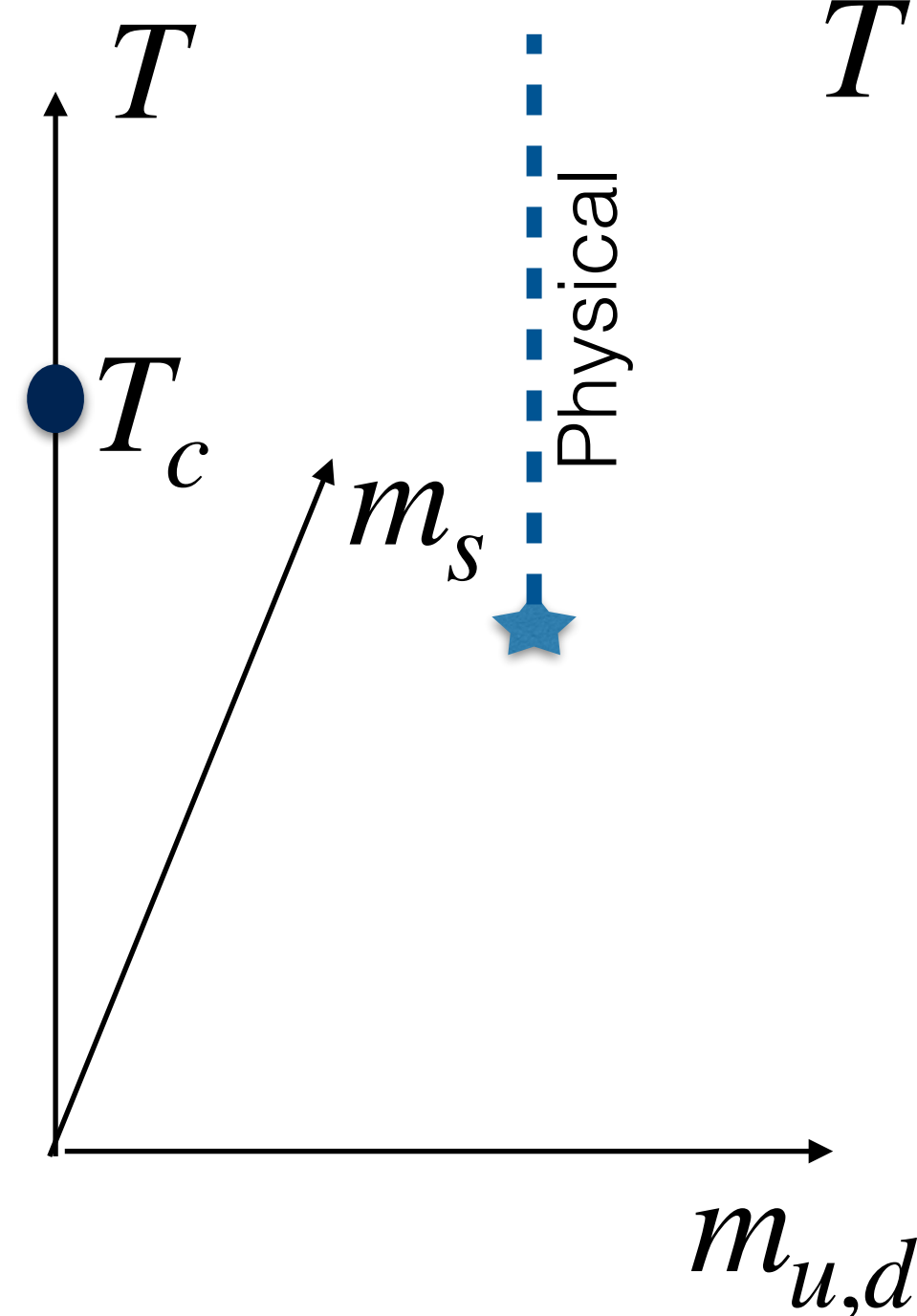


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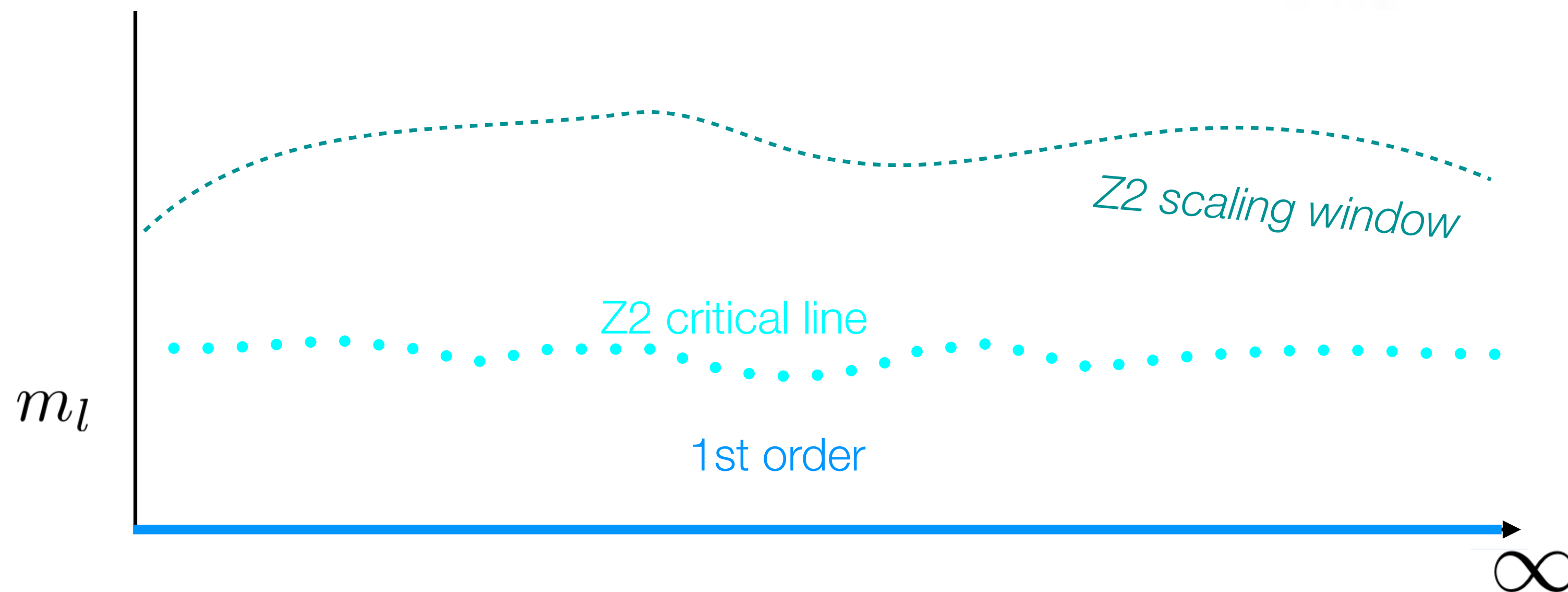
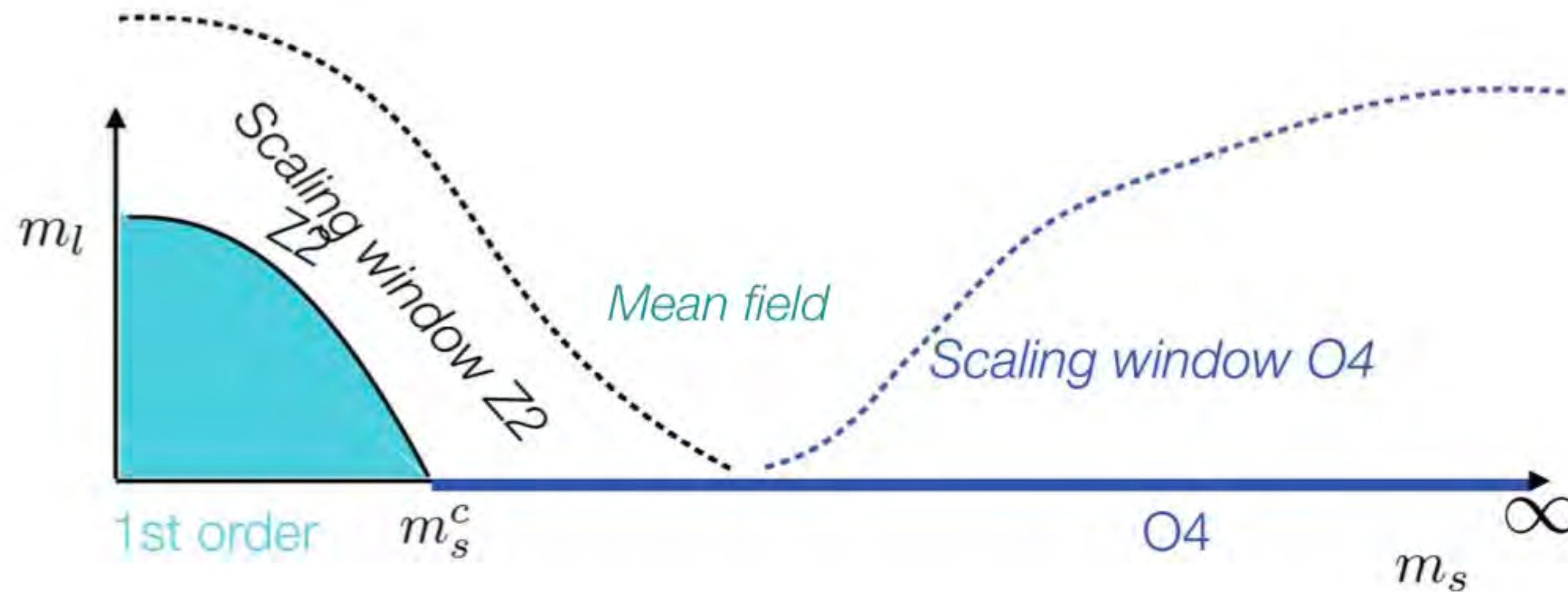
$SU_V(n)$



$T > T_c$ ($m = 0$): (which?) symmetry restoration $\Leftrightarrow$ order (universality)

$m \neq 0$: explicit symmetry breaking

$m_l \neq 0$, possible scenarios

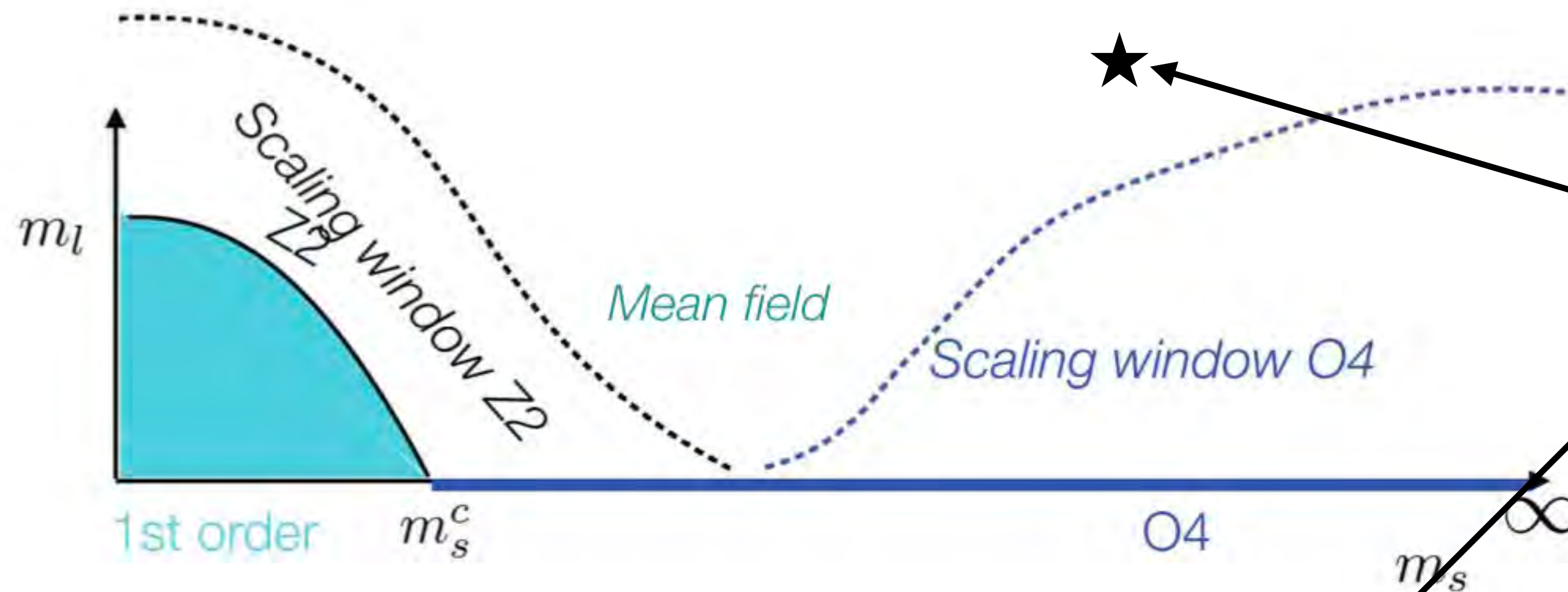


Scaling window:

universal behaviour
given by EoS

$$M = h^{1/\delta} f(t/h^{1/\beta\delta}) + \text{regular terms}$$

$m_l \neq 0$, possible scenarios

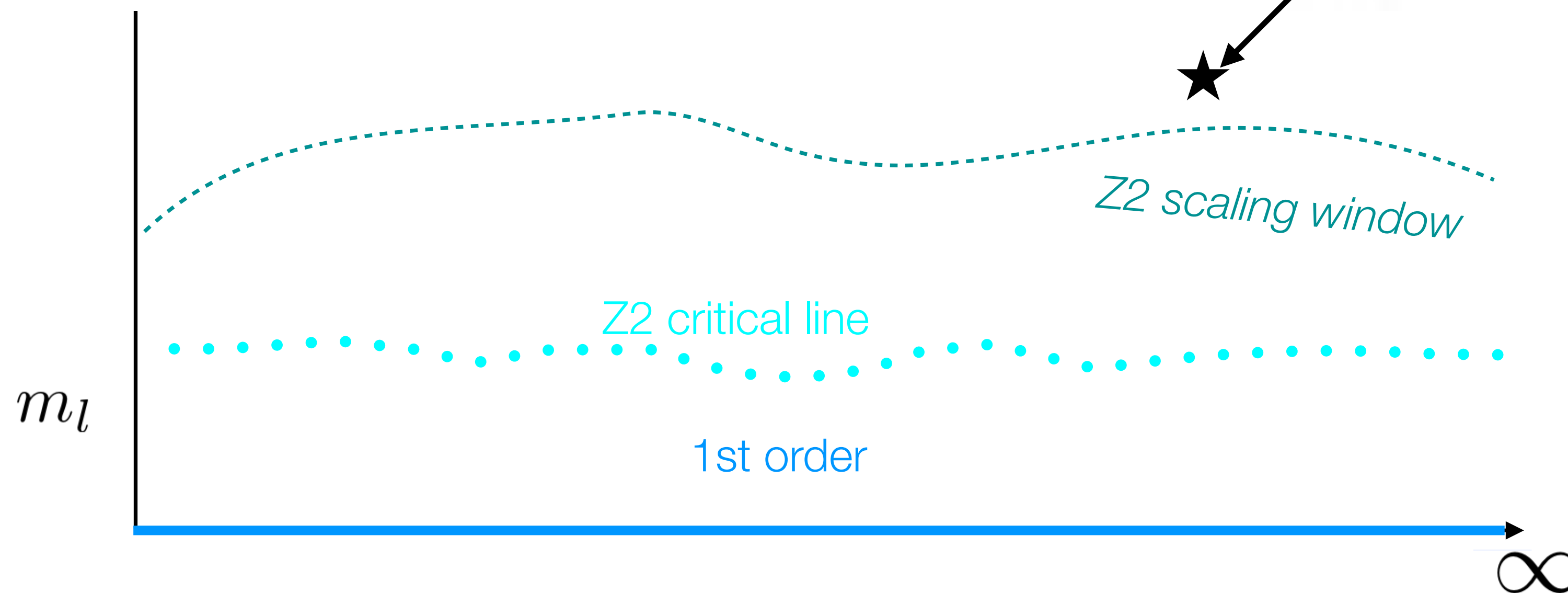


Physical point ?

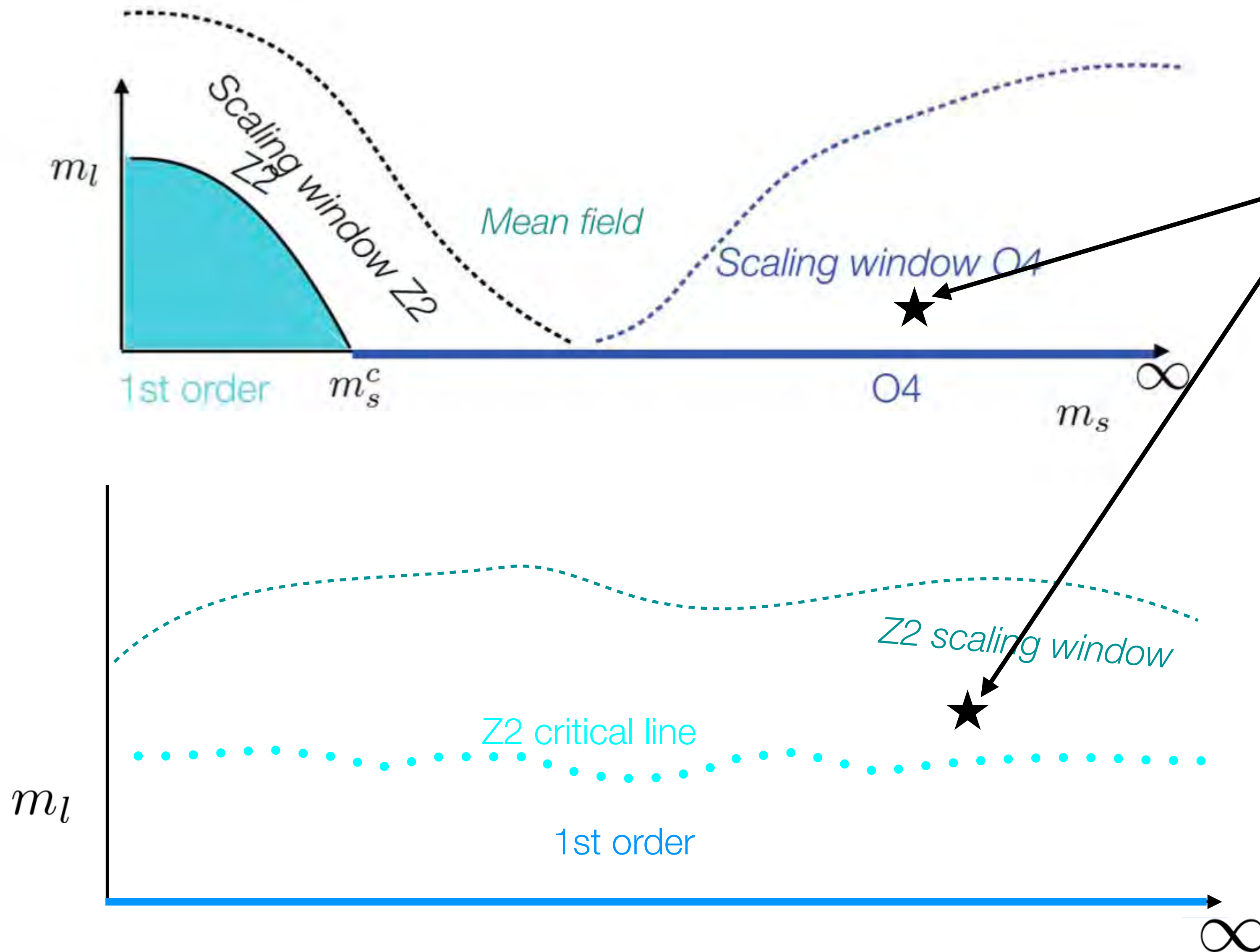
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$m_l \neq 0$, possible scenarios



Or here ?

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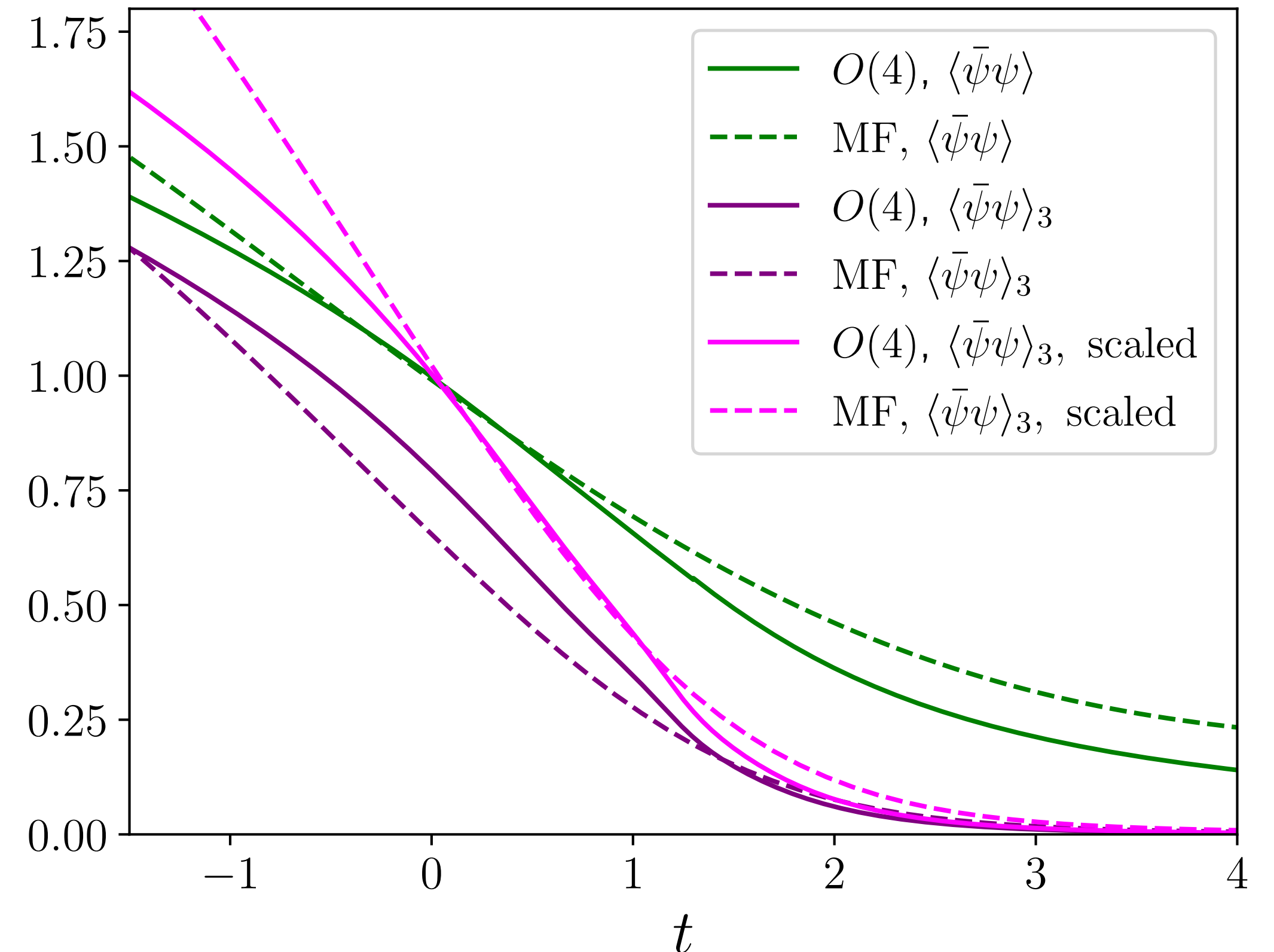
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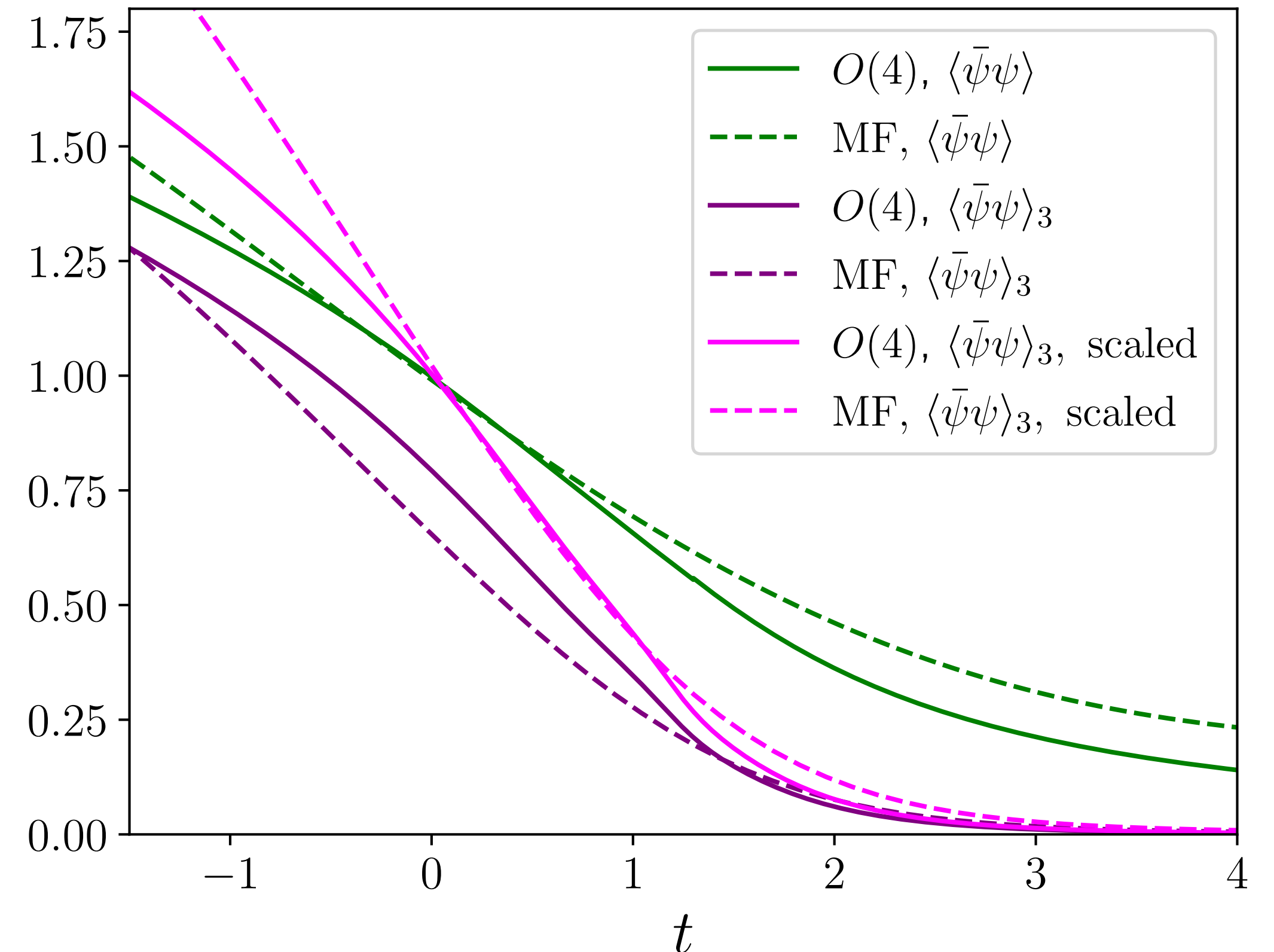
as $m \rightarrow 0$

- $\langle \bar{\psi}\psi \rangle_3 \sim t^{-\gamma-2\beta\delta}$ as $t \rightarrow \infty$
- $\langle \bar{\psi}\psi \rangle \sim t^{-\gamma}$ as $t \rightarrow \infty$

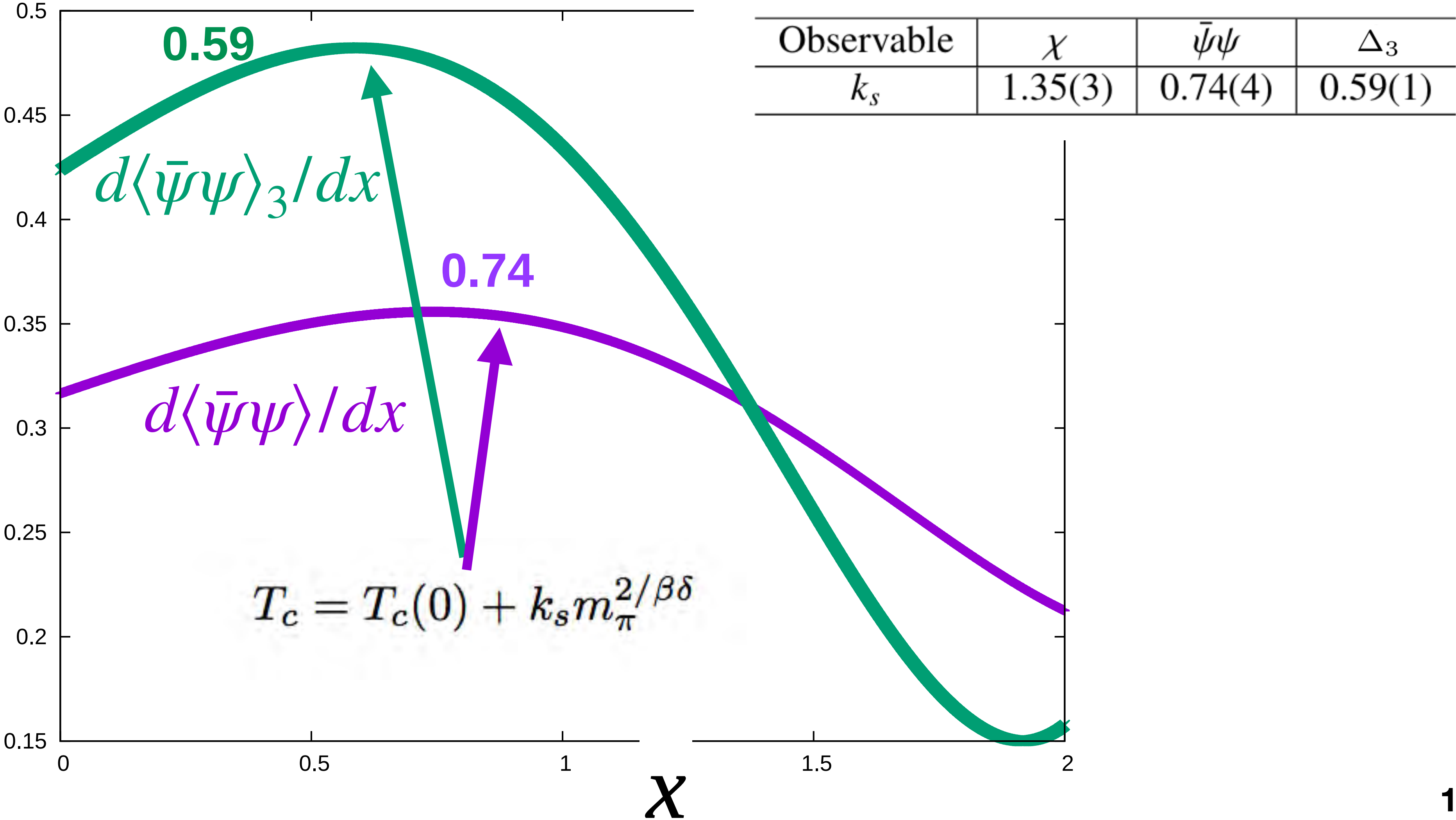


Novel order parameter

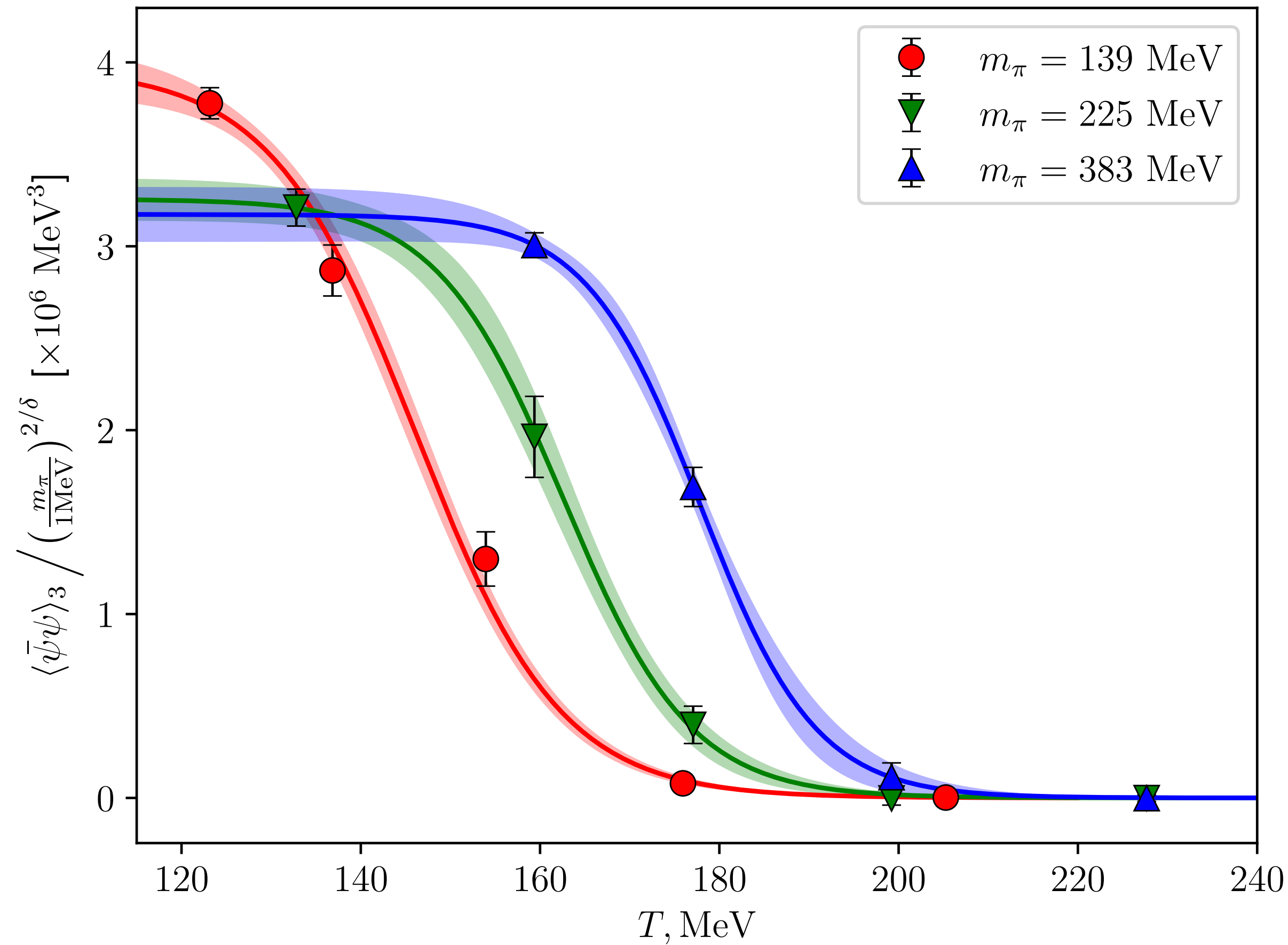
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as $m \rightarrow 0$
 - $\langle \bar{\psi}\psi \rangle_3 \sim t^{-\gamma-2\beta\delta}$ as $t \rightarrow \infty$
 - $\langle \bar{\psi}\psi \rangle \sim t^{-\gamma}$ as $t \rightarrow \infty$
- If $\langle \bar{\psi}\psi \rangle_3 < 0$: possibly first order, or
closeness to the first order phase transition
(Z_2 scenario ?)



Scaling of T_c with pion mass



Simple estimation of T_0 from EOS



Prediction of EoS:

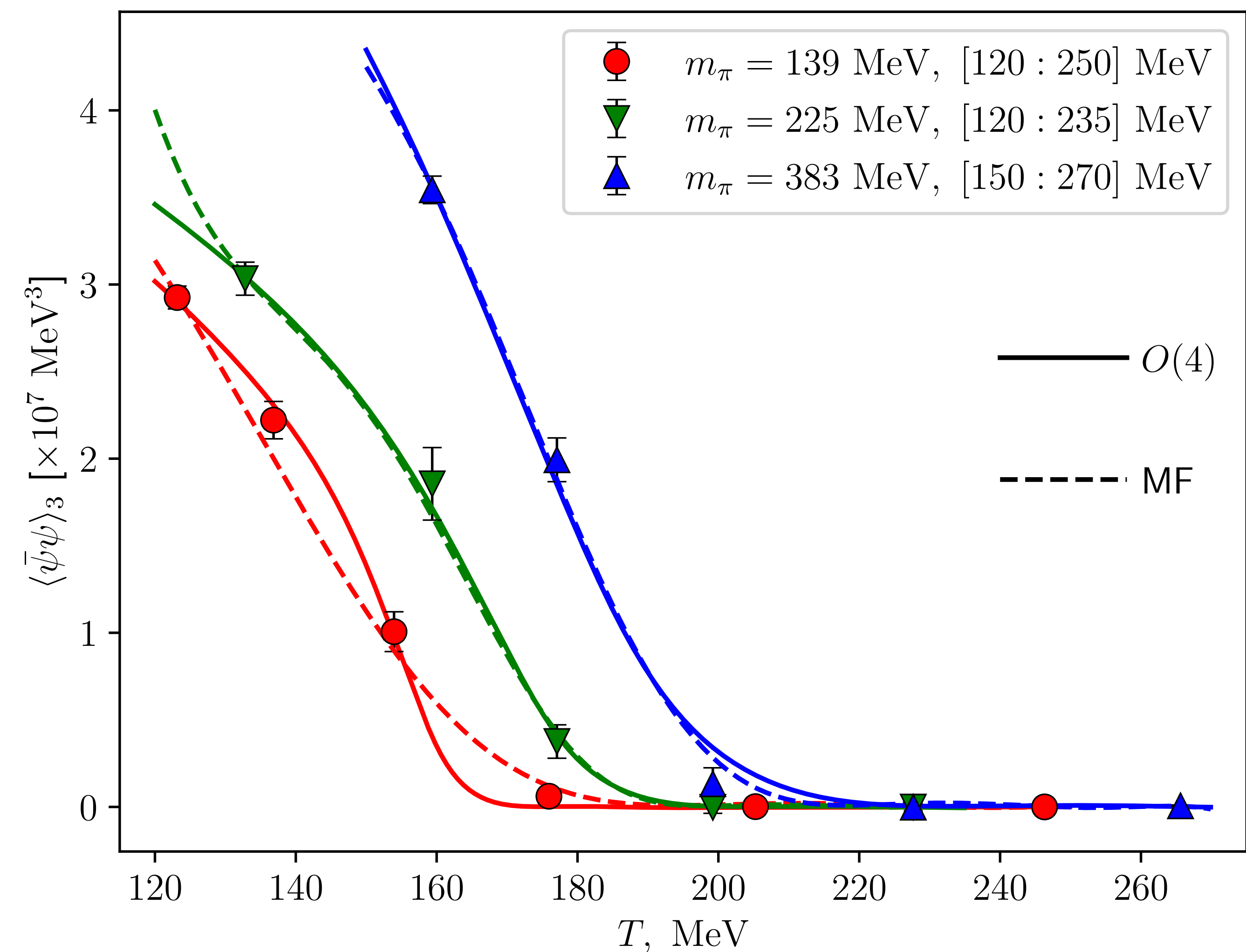
$$\frac{\langle \bar{\psi}\psi \rangle_3}{m^{1/\delta}} \sim \frac{\langle \bar{\psi}\psi \rangle_3}{m_\pi^{2/\delta}} = \text{const}$$

at

$$T = T_0(m_\pi = 0) = 138(2) \text{ MeV}$$

$$M = h^{1/\delta} f(t/h^{1/\beta\delta}) + \text{regular terms}$$

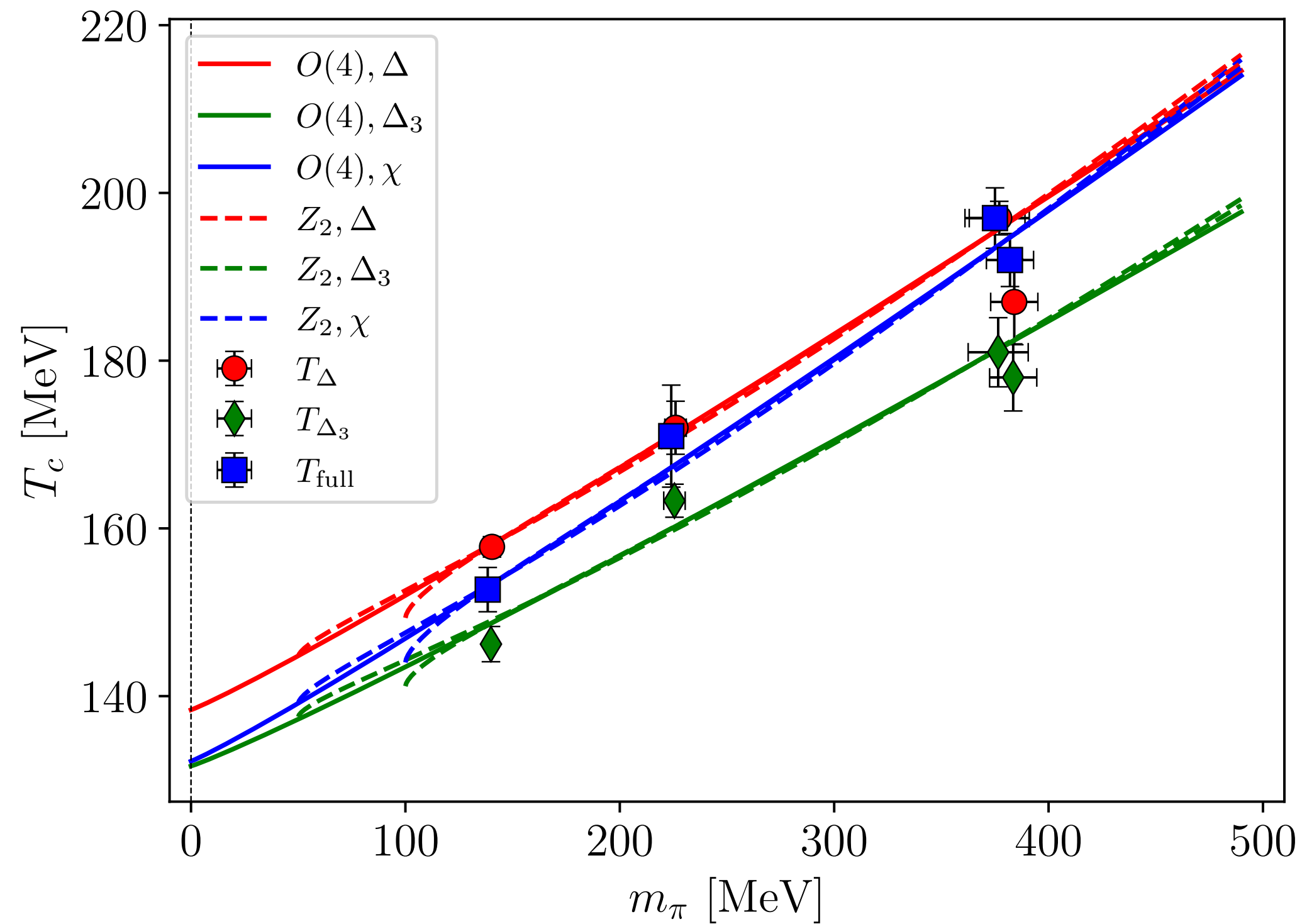
O(4) vs mean field



Mild tension between data and MF for $m_\pi=139$ MeV

m_π [MeV]	T_0 [MeV]
139	142(2)
225	159(3)
383	174(2)

Z₂ vs O(4) scaling



$$T_0 = T_c(m_\pi \rightarrow 0) = 134^{+6}_{-4} \text{ MeV}$$

O(4) scaling:

Observable	T_0 [MeV]	$z_p/z_{\bar{\psi}\psi_3}$	$z_p/z_{\bar{\psi}\psi_3} O(4)$	$z_p O(4)$
χ	132(4)	1.24(17)	2.45(4)	1.35(3)
$\langle\bar{\psi}\psi\rangle$	138(2)	1.15(24)	1.35(7)	0.74(4)
$\langle\bar{\psi}\psi\rangle_3$	132(3)	1	1	0.55(1)

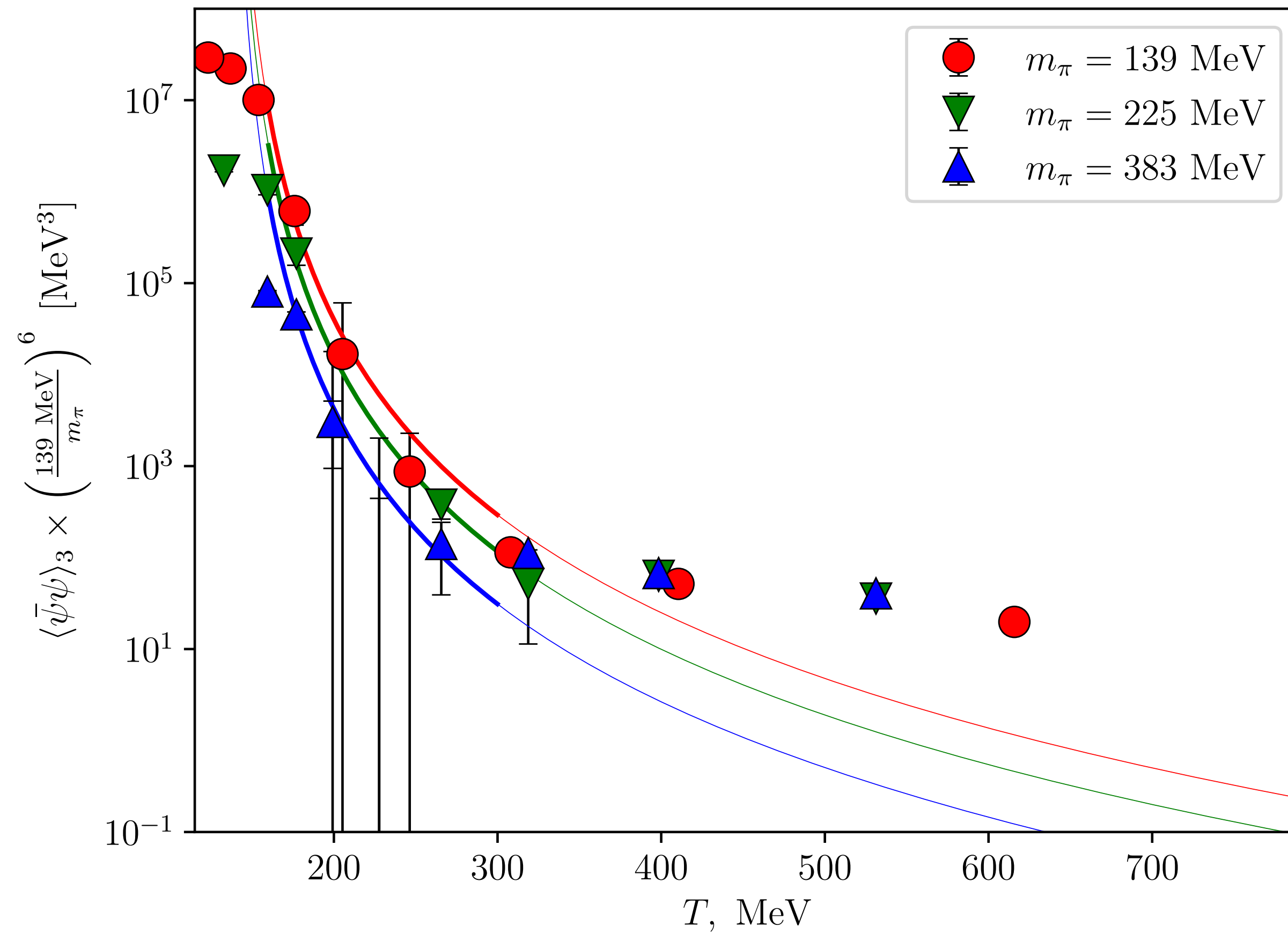
Z₂ scaling:

$m_\pi^c = 100 \text{ MeV}$ is still ok

$m_\pi^c = 0 \text{ MeV}$ is indistinguishable from O(4)

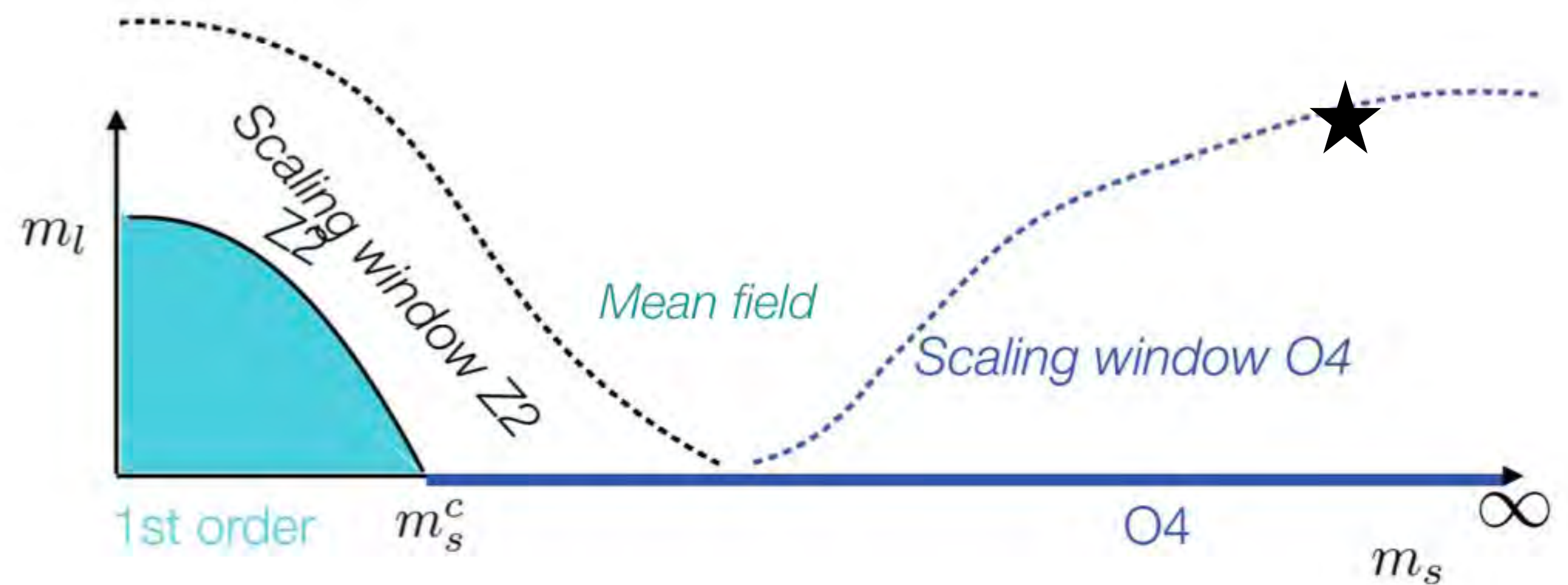
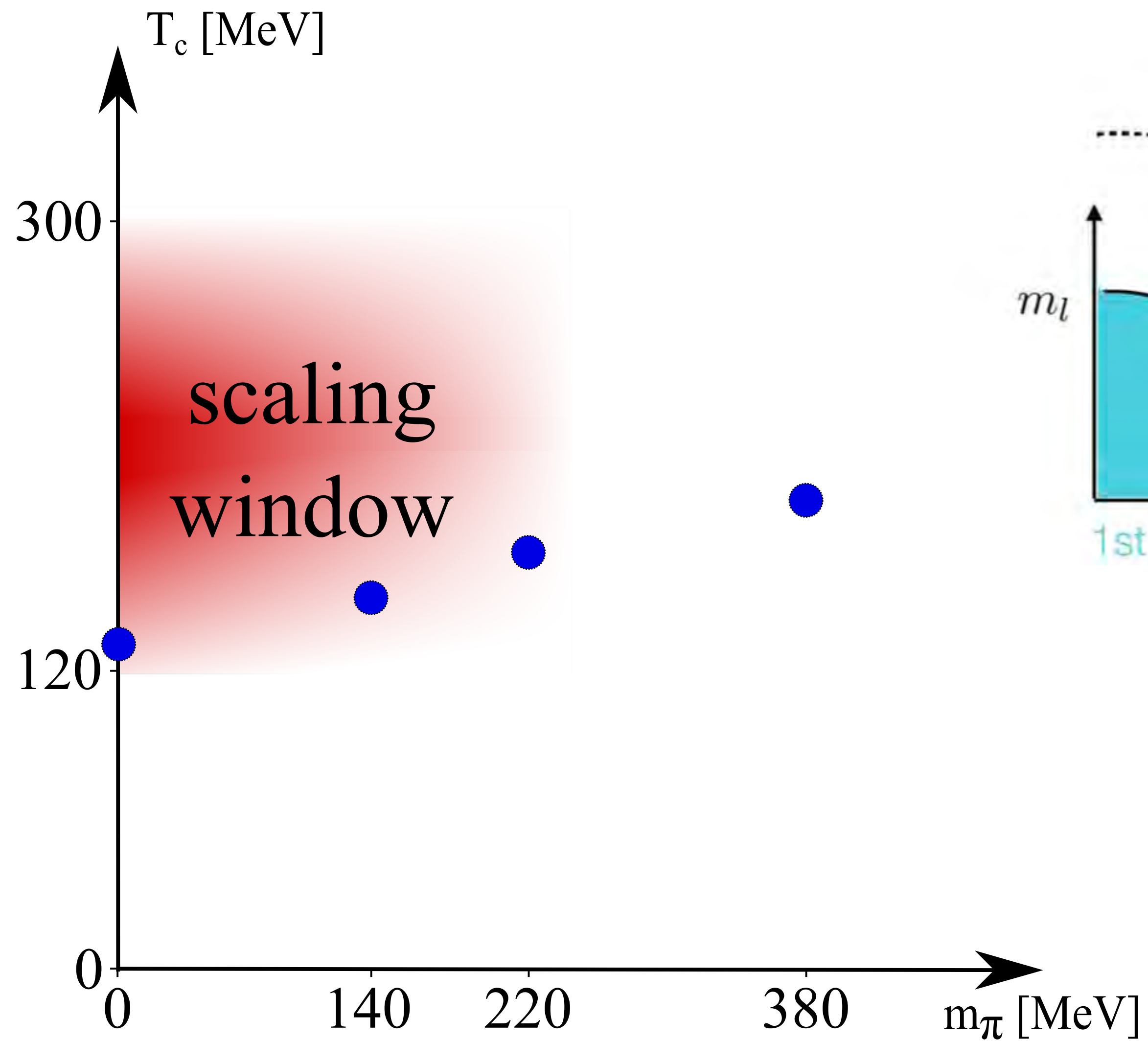
$$T_c = T_c(0) + k_s(m - m_c)^{1/\beta\delta}, m \sim m_\pi^2$$

Large temperature behaviour



- O(4): $\langle \bar{\psi}\psi \rangle_3 \sim t^{-\gamma-2\beta\delta}$
- Griffith analyticity:
 $\langle \bar{\psi}\psi \rangle_3 \sim m^3 \sim m_\pi^6$
- $T \sim 300 \text{ MeV}$

Scaling window



Threshold at $T = 300$ MeV and topology

Method to measure topological susceptibility

QCD and topology, finite temperature

Method to measure topological susceptibility

QCD and topology, finite temperature

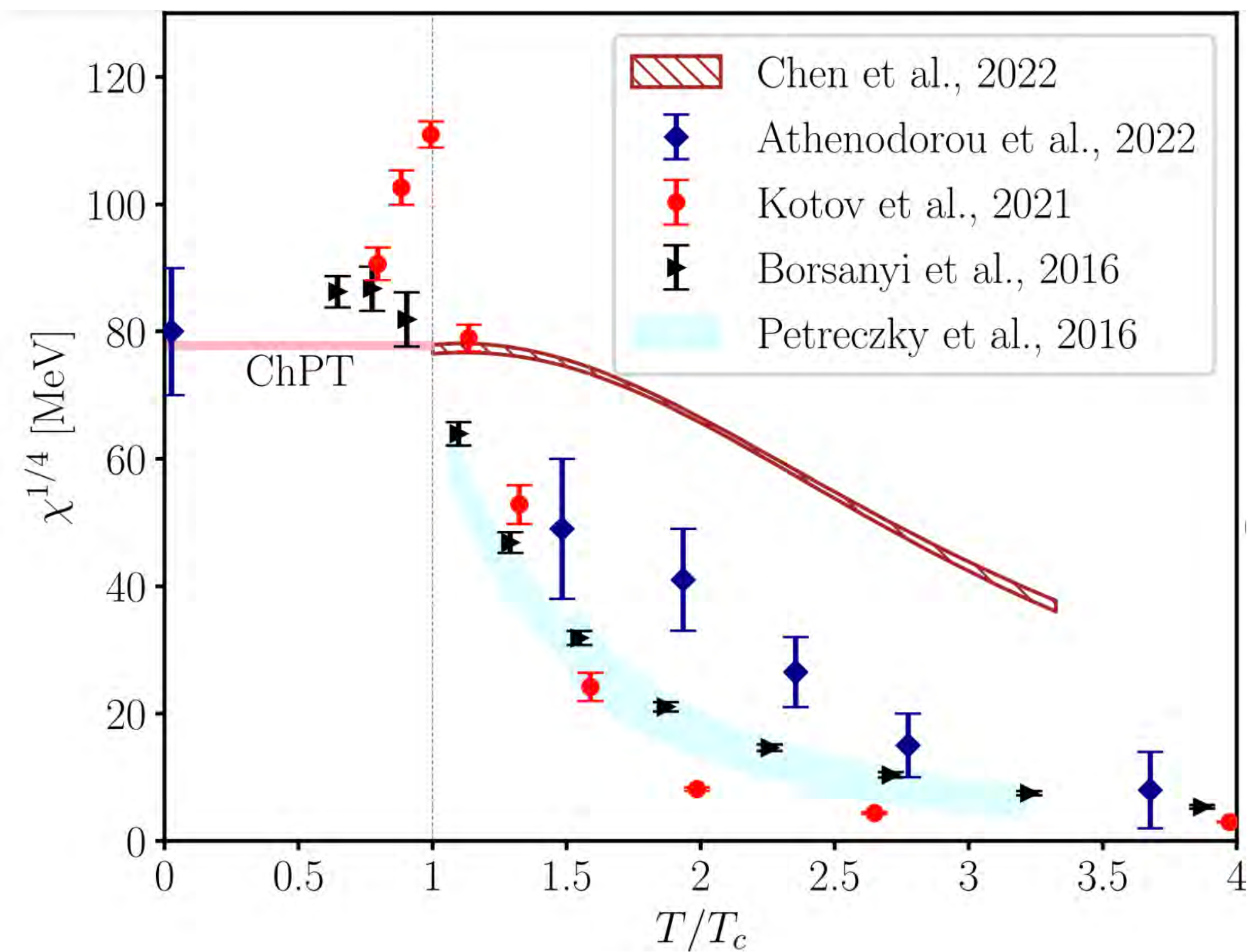
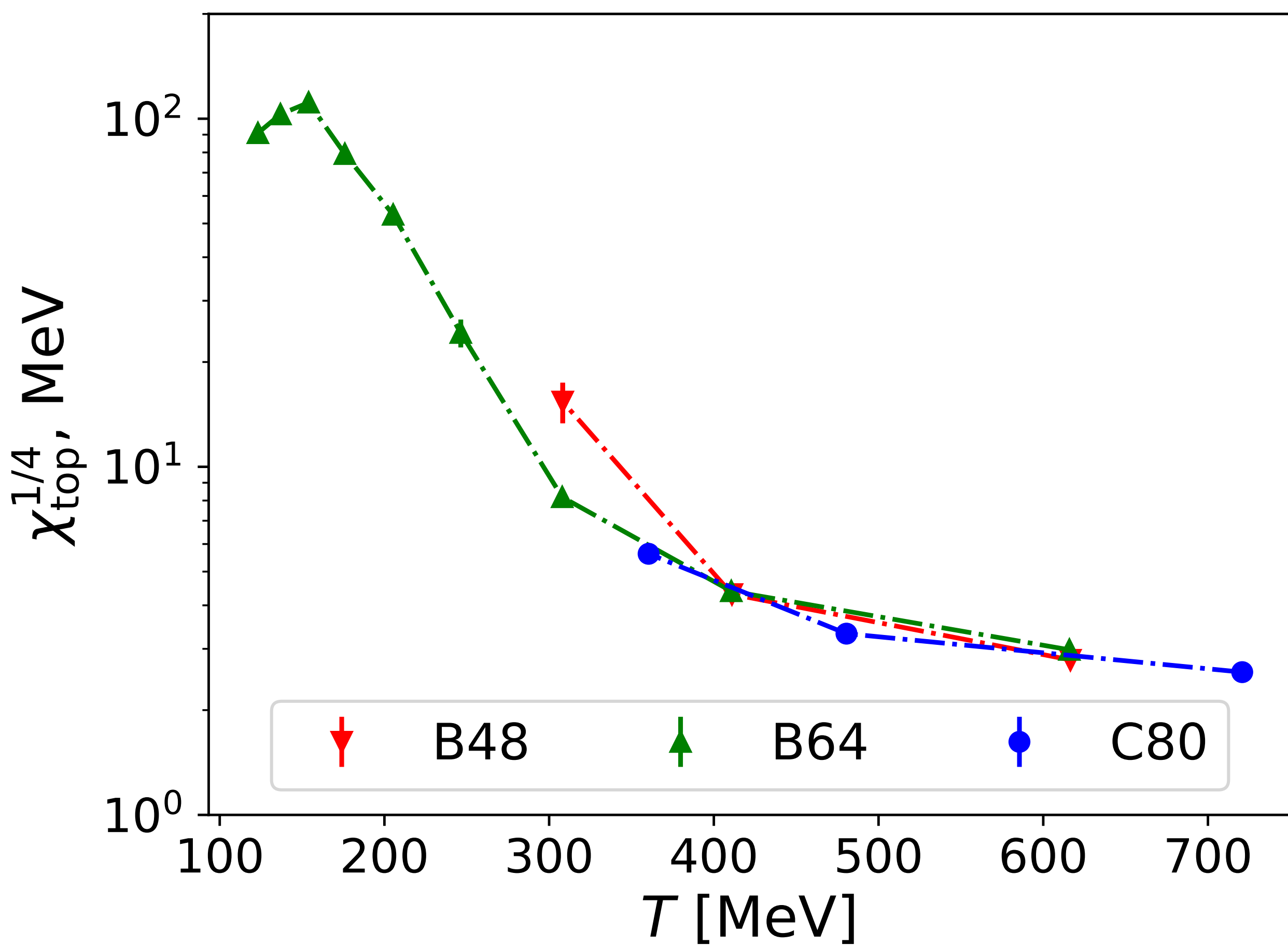
$$\chi_{\text{top}} = \langle Q_{\text{top}} \rangle^2 / V = m_l^2 \chi_{5,\text{disc}}$$

$$m \int d^4x \bar{\psi} \gamma_5 \psi = Q_{\text{top}}$$

[Kogut, Lagae, Sinclair, 1998]

$$\begin{array}{ccc}
 & SU_L(2) \times SU_R(2) & \\
 \chi_{5,\text{con}} & \pi : \bar{\psi} \gamma_5 \frac{\tau}{2} \psi & \longleftrightarrow \sigma : \bar{\psi} \psi \quad \chi_{\text{con}} + \chi_{\text{disc}} \\
 & \uparrow U_A(1) & \uparrow U_A(1) \\
 \chi_{\text{con}} & \delta : \bar{\psi} \frac{\tau}{2} \psi & \longleftrightarrow \eta : \bar{\psi} \gamma_5 \psi \quad \chi_{5,\text{con}} - \chi_{5,\text{disc}} \\
 & SU_L(2) \times SU_R(2) &
 \end{array}$$

$$\chi_{\pi} - \chi_{\delta} = \chi_{\text{disc}} = \chi_{5,\text{disc}}, \quad \text{for } T \geq T_C, m_l \rightarrow 0 \quad \Rightarrow \chi_{\text{top}} = \langle Q_{\text{top}} \rangle^2 / V = m_l^2 \chi_{\text{disc}}$$

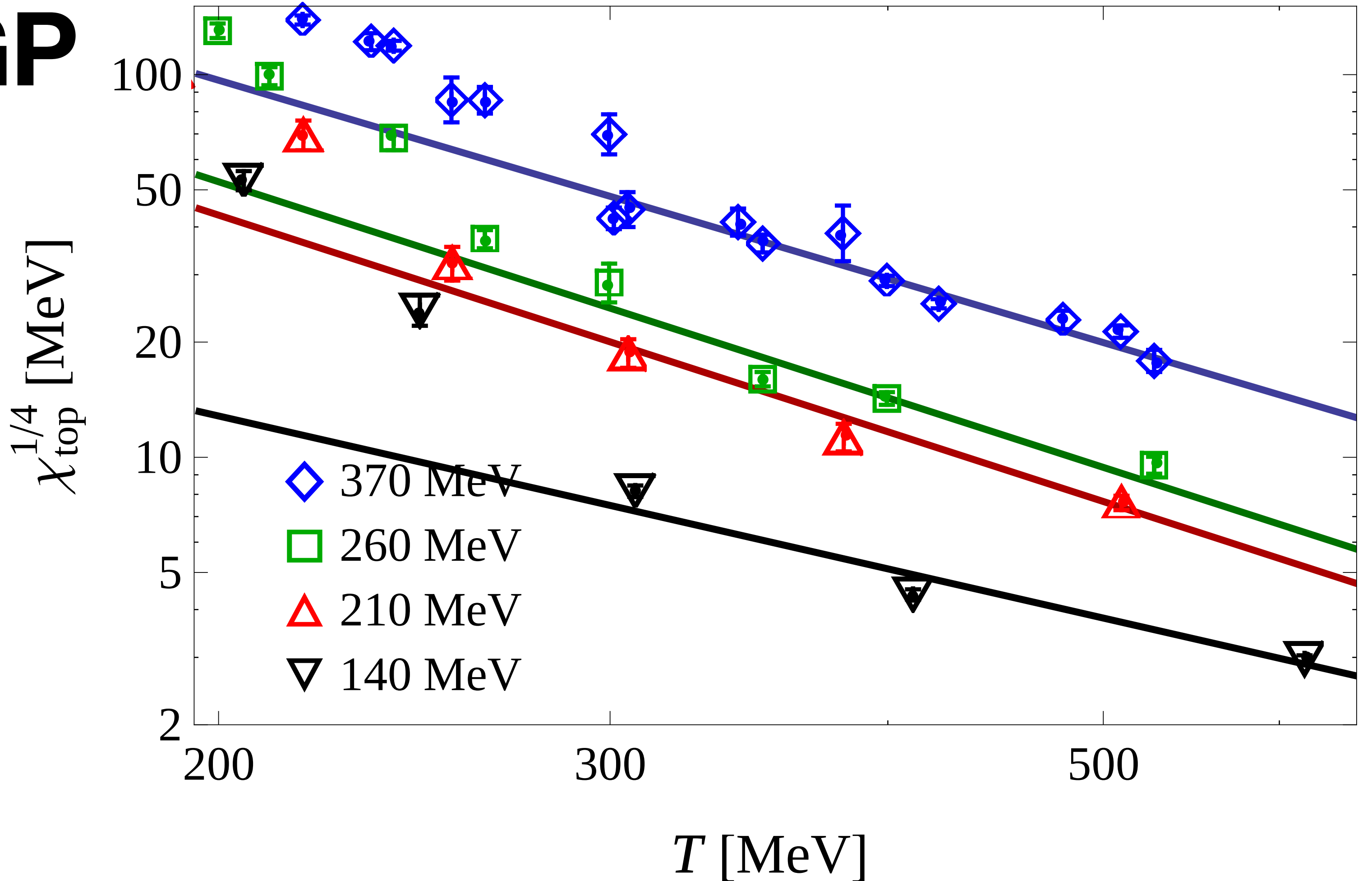


Threshold in QGP

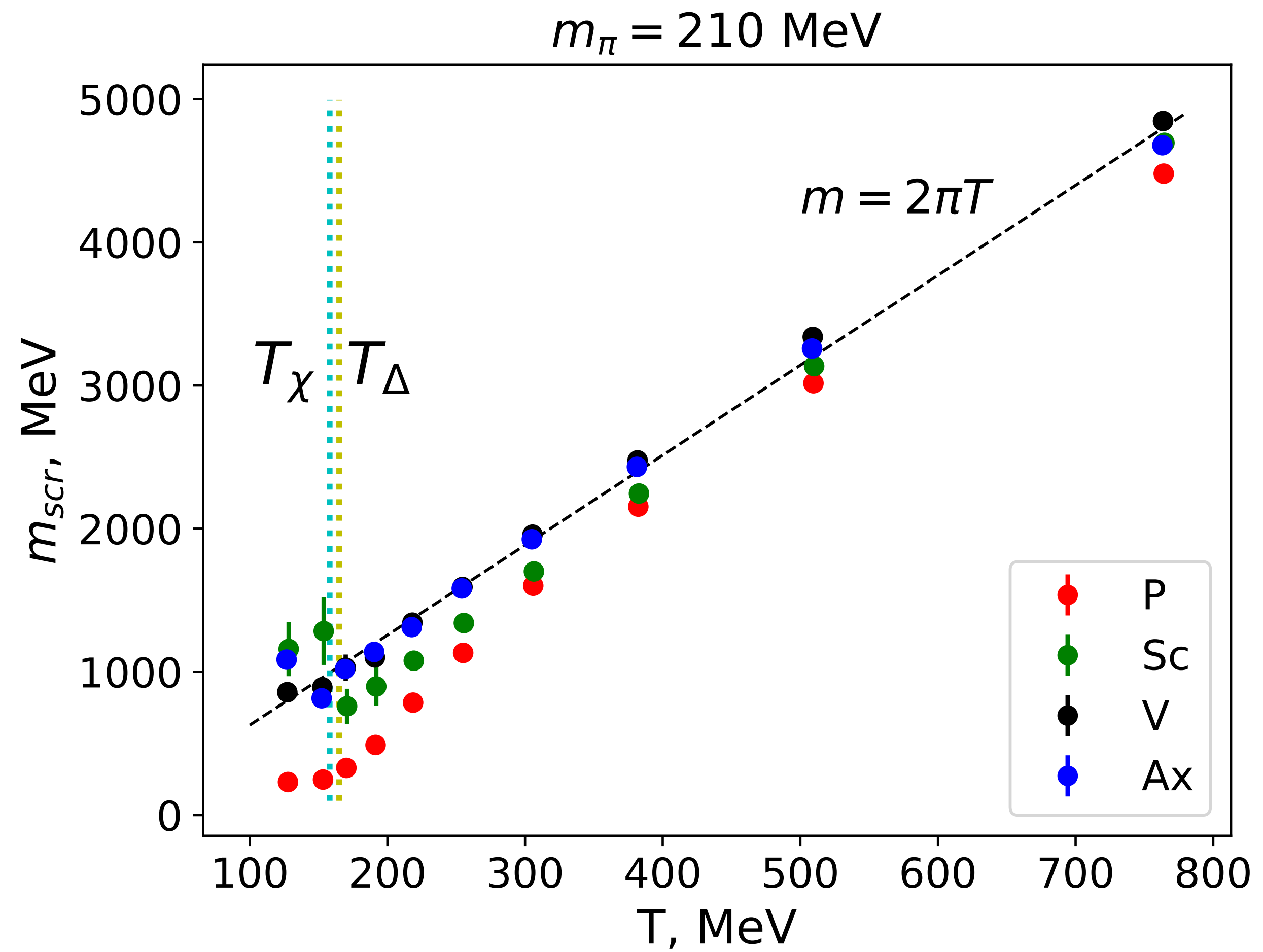
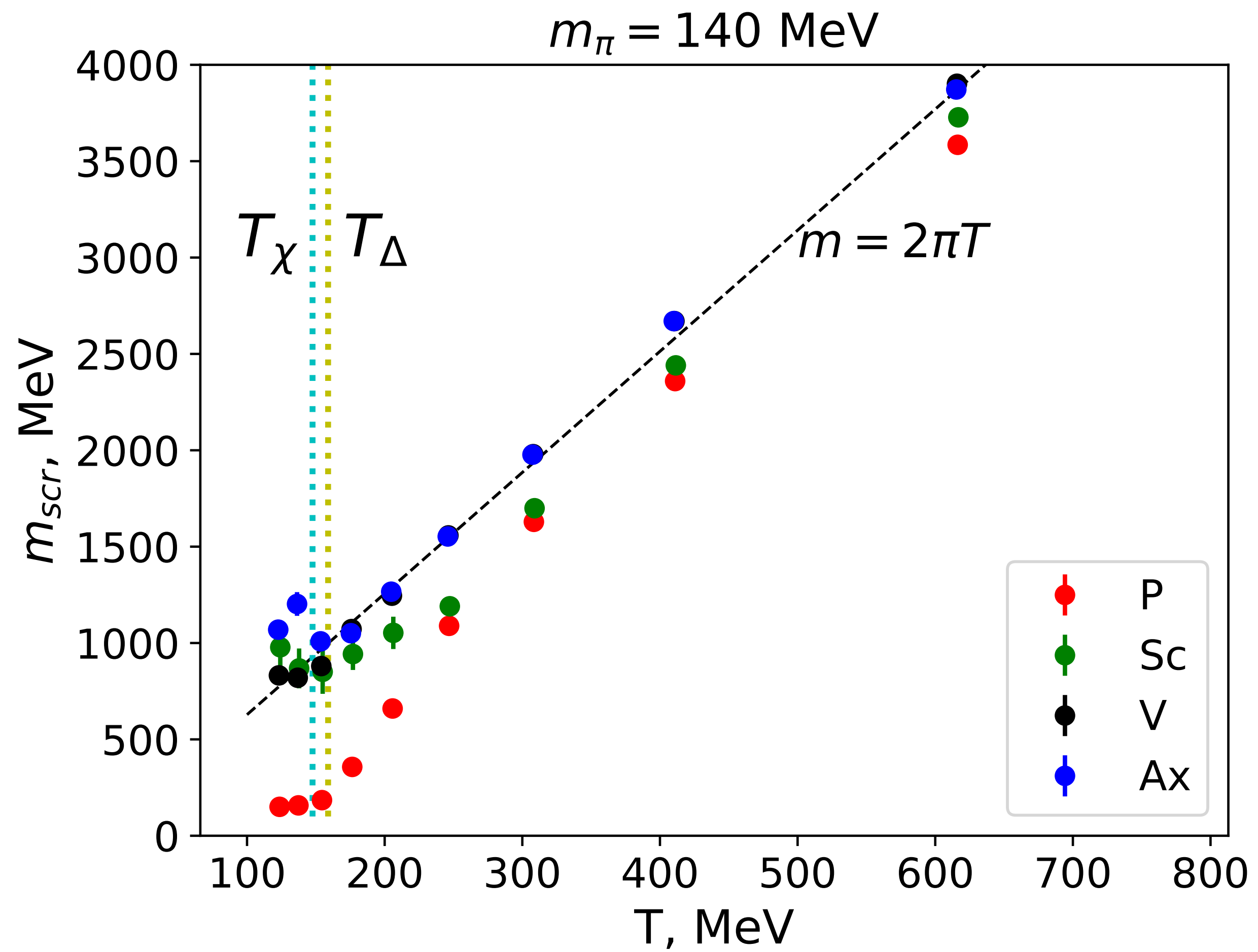
$T \sim 300 \text{ MeV}$

- Onset of DIGA behaviour

$$\chi \sim T^{-d}$$



Screening masses



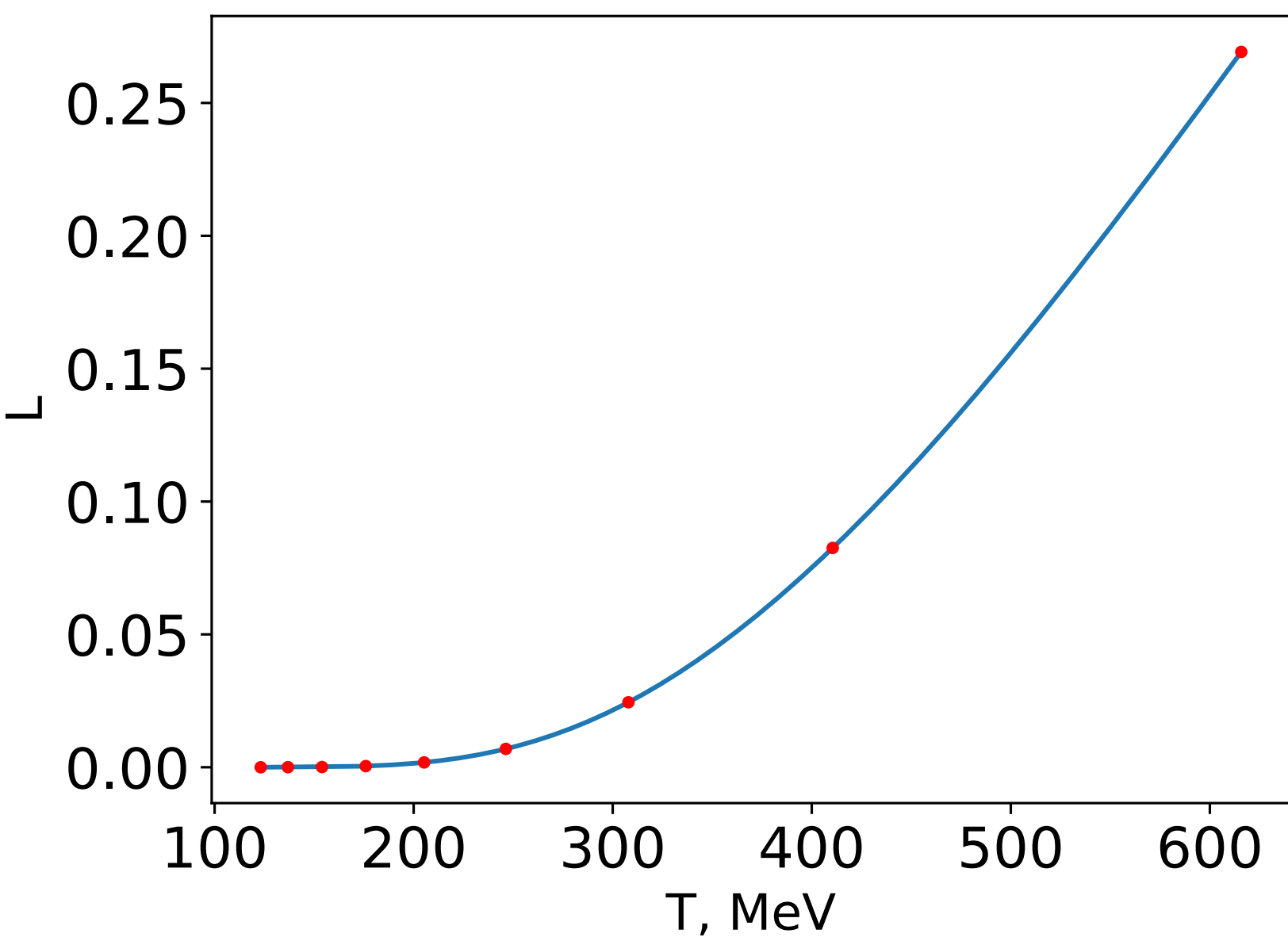
Polyakov loop

- Polyakov loop $|\langle P \rangle| = e^{-F_Q/T}$: needs renormalisation
- Renormalise with Gradient Flow
- GF evolution: $\partial_\tau A_\mu(\tau, t, x) = -\frac{\partial S}{\partial A_\mu}$, smearing over radius $f = \sqrt{8\tau}$
- Different GF times $\implies$ different renormalisation schemes:

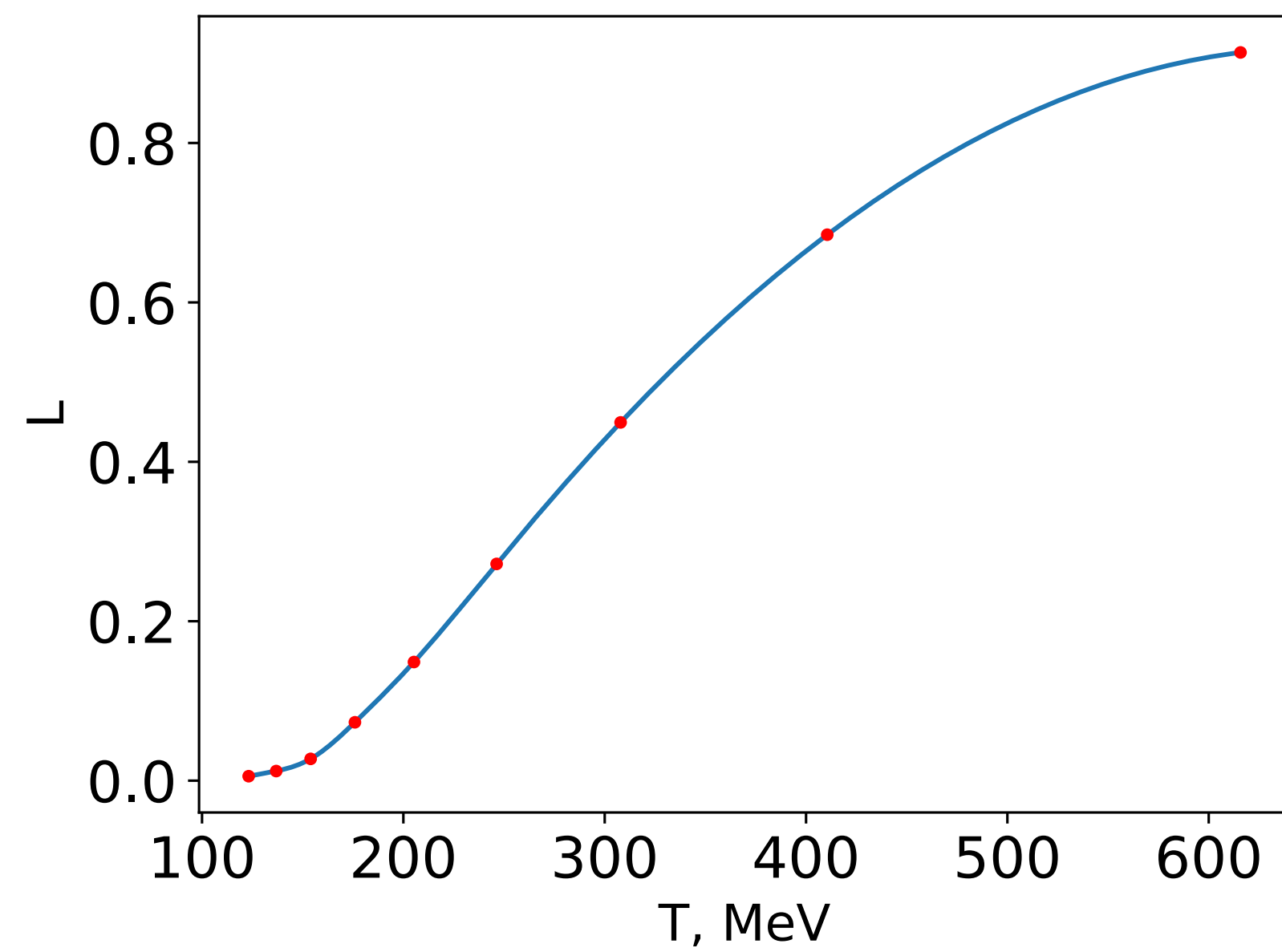
$$F_Q(\tau', T) = F_Q(\tau, T) + \delta F_Q(\tau, \tau')$$
- Inflection point of Polyakov loop is ambiguously defined
- Use $F_Q(T)$ or $S_Q(T) = -\partial F_Q(T)/\partial T$

[TUMQCD, 2016]

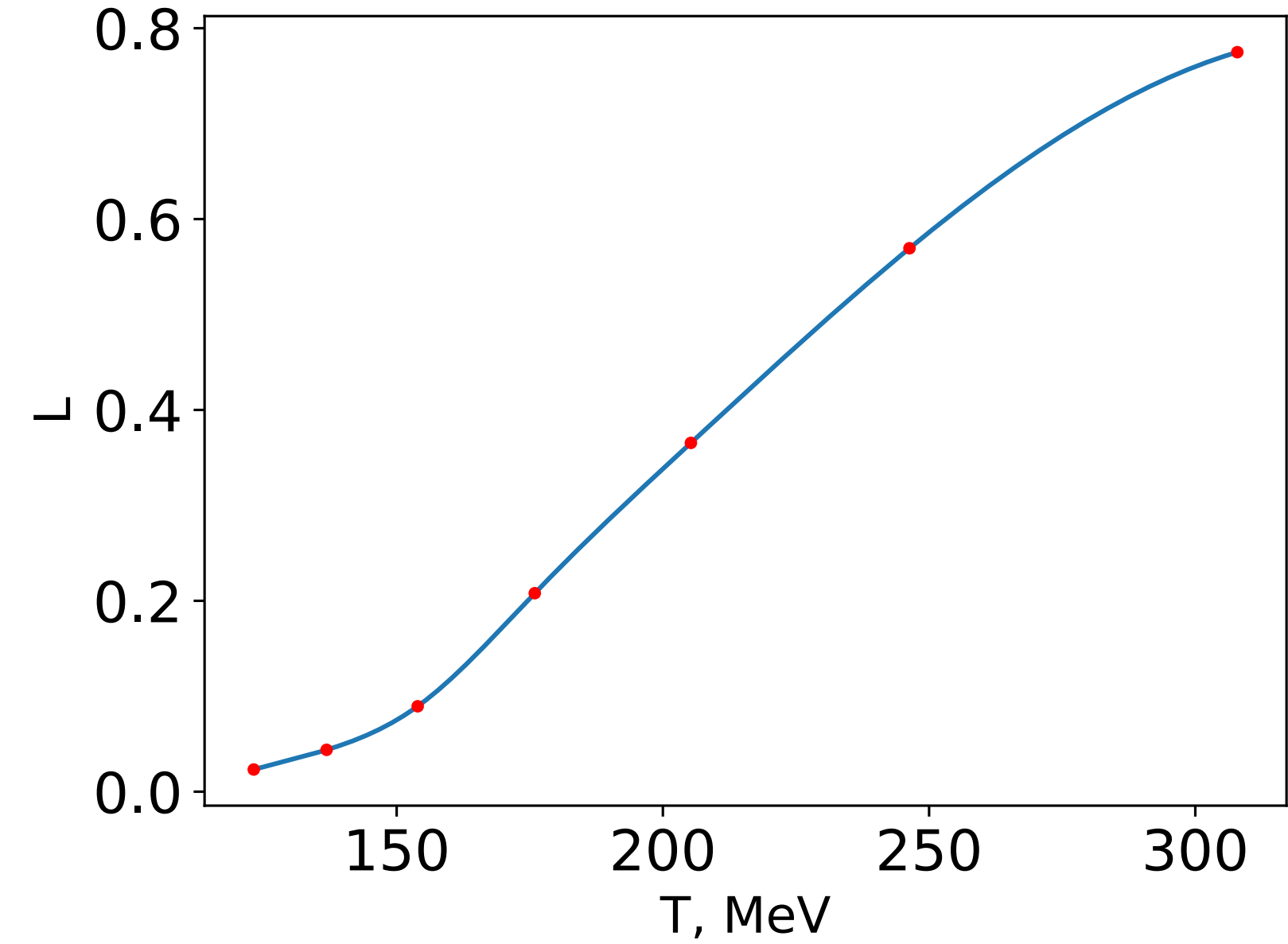
$f = 0.000$ fm

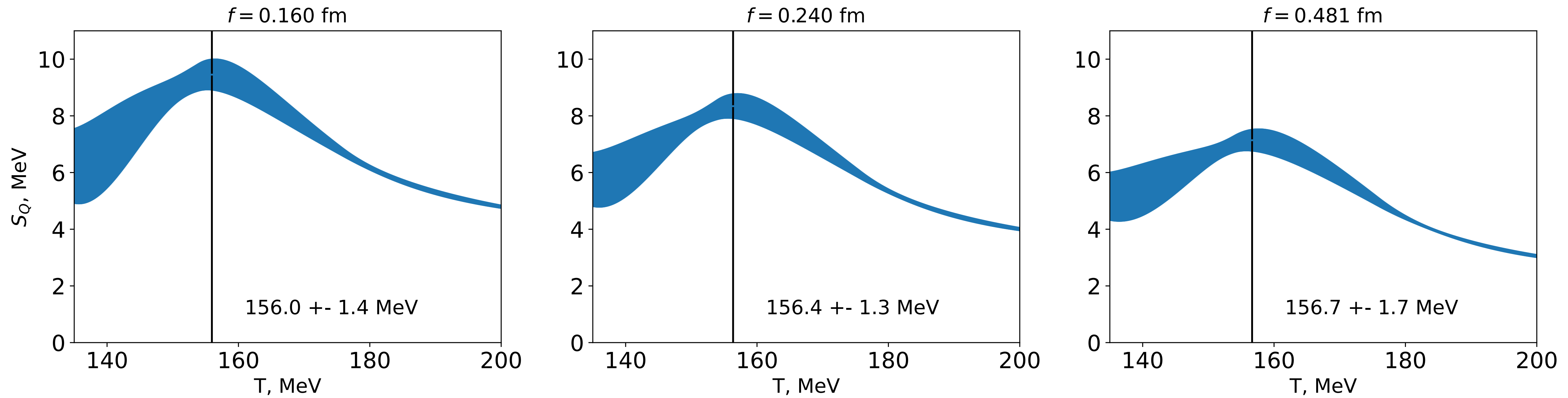


$f = 0.240$ fm



$f = 0.481$ fm





Consistent with chiral pseudo-critical temperature

Conclusions

- Thermal QCD phases by TWEXT collaboration
- $\langle \bar{\psi}\psi \rangle_3 = \langle \bar{\psi}\psi \rangle - m\chi$ for critical/scaling phenomena
- $T_0 = 134^{+6}_{-4}$ MeV in the chiral limit $m \rightarrow 0$
- Consistent with $O(4)$ scaling for $m_\pi \lesssim 140$ MeV, $T \in [120, 300]$ MeV
- $T \sim 300$ MeV: indications of threshold in QGP

