

Making sense of Green's function integrals

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Analytic and Algebraic Methods in Physics XXII
Czech Technical University, Prague, August 2025

This research was funded by the Austrian Science Fund (FWF)

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- The Green's function integral

$$\psi(t, x) = \int_{\mathbb{R}} G(t, x, y) F(y) dy$$

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$$\int_{\mathbb{R}} e^{iax^2} f(x) dx.$$

- Method I: Fresnel integral
- Method II: Integration by parts
- Application: Superoscillations

The Schrödinger equation

For some potential V and initial conditions F , we consider

$$\begin{aligned} i \frac{\partial}{\partial t} \psi(t, x) &= \left(- \frac{\partial^2}{\partial x^2} + V(t, x) \right) \psi(t, x), & t > 0, x \in \mathbb{R}, \\ \psi(0, x) &= F(x), & x \in \mathbb{R}. \end{aligned}$$

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We want to solve it using the **Green's function integral**

$$\psi(t, x) = \int_{\mathbb{R}} G(t, x, y) F(y) dy.$$

What kind of integrals are we faced with?

$$\Psi(t, x) = \int_{\mathbb{R}} \underbrace{G(t, x, y)}_{\substack{\text{bounded,} \\ \text{but of special shape}}} \underbrace{F(y)}_{\text{bounded}} dy.$$

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Pöschl-Teller $V \sim \frac{1}{\cosh^2(x)}$

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$$G(t, x, y) = \frac{\Theta(xy)\sqrt{xy}}{2i^{\nu+1}t} J_{\nu}\left(\frac{xy}{2t}\right) e^{i\frac{x^2+y^2}{4t}} \quad \text{Centrifugal potential } V \sim \frac{1}{x^2}$$

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$$G(t, x, y) = -\frac{c}{2}\Lambda\left(\frac{|x|+|y|}{2\sqrt{it}}\right) + c\sqrt{it} e^{i\frac{(|x|+|y|)^2}{4t}} \quad \delta\text{-potential } V \sim \delta(x)$$

Oscillatory integrals

In all the cases we are faced with **oscillatory integrals** of the form

$$\int_{\mathbb{R}} e^{iy^2} f(y) dy, \quad (1)$$

for functions $f \notin L^1(\mathbb{R})$.

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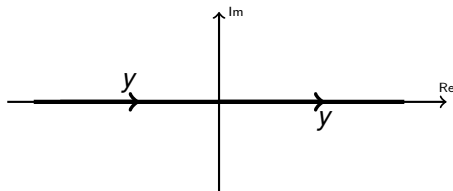
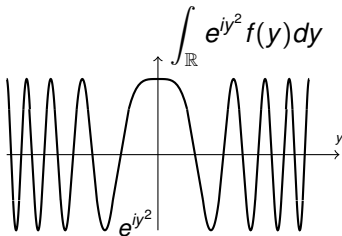
- Two **methods** (with variations) to give meaning to (1).
- Different **assumptions on the function f** .
- Find **absolute convergent representations** of (1).

Method I:

Fresnel integral

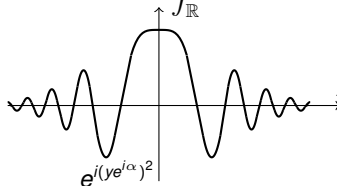
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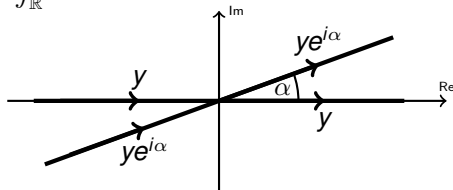
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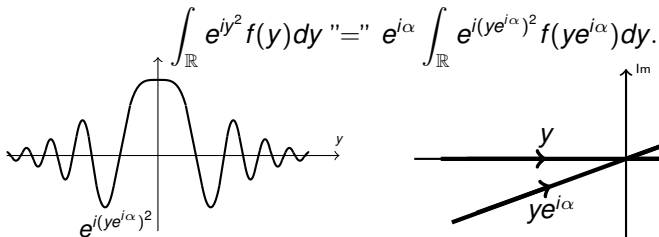
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$$\int_{\mathbb{R}} e^{iy^2} f(y) dy \stackrel{""}{=} e^{i\alpha} \int_{\mathbb{R}} e^{i(ye^{i\alpha})^2} f(ye^{i\alpha}) dy.$$




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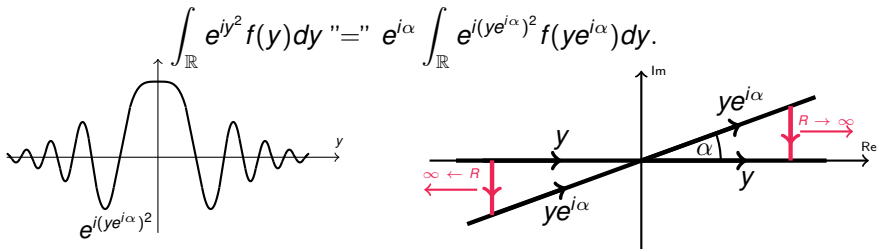
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Formally, this is the **substitution** $y \rightarrow ye^{i\alpha}$.

Fresnel integral

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Strictly speaking this is the **Cauchy theorem** along a triangle path

$$\begin{aligned} \int_{-R}^R e^{iy^2} f(y) dy &= e^{i\alpha} \int_{-R}^R e^{i(ye^{i\alpha})^2} f(ye^{i\alpha}) dy \\ &+ \int_{-R}^{-R-iR \tan \alpha} e^{iz^2} f(z) dz + \int_R^{R+iR \tan \alpha} e^{iz^2} f(z) dz \end{aligned}$$

Assumptions on f

We need the vertical **red paths** to vanish when $R \rightarrow \infty$!

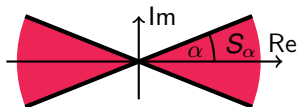
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- The function f **extends holomorphically** to the double sector S_α .

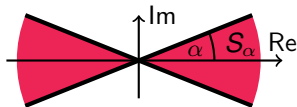


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- The function f is **exponentially bounded** by

$$|f(z)| \leq A e^{B|\operatorname{Im}(z)|}, \quad z \in S_\alpha.$$

Then we obtain the convergence

$$\left| \int_R^{R+iR \tan \alpha} e^{iz^2} f(z) dz \right| \leq A \int_0^{R \tan \alpha} e^{-2Ry} e^{By} dy \leq \frac{A}{2R-B} \xrightarrow{R \rightarrow \infty} 0.$$

Definition of the Oscillatory integral

Definition 1a) of Oscillatory integral

If f is holomorphic on S_α and bounded by $|f(z)| \leq Ae^{B|\operatorname{Im}(z)|}$, then

$$\int_{\mathbb{R}} e^{iy^2} f(y) dy := \lim_{R \rightarrow \infty} \int_{-R}^R e^{iy^2} f(y) = \lim_{R \rightarrow \infty} e^{i\alpha} \int_{-R}^R e^{i(ye^{i\alpha})^2} f(ye^{i\alpha}) dy$$

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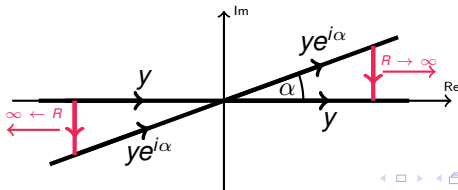
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But we can do better with a **Gaussian regularizer**, i.e. for some $\varepsilon > 0$

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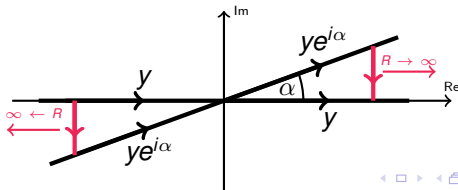
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If we now assume

$$|f(z)| \leq Ae^{B|z|}, \quad z \in S_\alpha,$$

there vanishes the terms

$$\left| \int_R^{R+iR \tan \alpha} e^{-\varepsilon z^2} e^{iz^2} f(z) dz \right| \leq Ae^{-\varepsilon R^2} \int_0^R e^{\varepsilon y^2 - 2Ry} e^{B|R+iy|} dy \xrightarrow{R \rightarrow \infty} 0.$$

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Y. Aharonov, J. Behrndt, F. Colombo, P. S., A unified approach to Schrödinger evolution of superoscillations and supershifts. J. Evol. Eq. **22** (2022).



P. S., Time evolution of superoscillations for the Schrödinger equation on $\mathbb{R} \setminus \{0\}$, Quantum Stud. Math. Found. **9** (2022).

Method II:

Integration by parts

The basic formula

The main idea is to use the **integration by parts formula**

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$$\int_{\mathbb{R}} e^{iy^2} f(y) dy = \frac{1}{2i} \int_{\mathbb{R}} \frac{d}{dy} (e^{iy^2}) \frac{f(y)}{y} dy$$

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$$\begin{array}{ccc} f & \longrightarrow & \frac{f'}{y} \\ \downarrow & & \\ \frac{f}{y^2} & & \end{array}$$

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The following diagram visualizes this method.

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Instead of f , we have to integrate the functions $\frac{f'}{y}$ and $\frac{f}{y^2}$, with the additional **decaying factors** $\frac{1}{y}$ and $\frac{1}{y^2}$.

Inductively applying this formula

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Inductively applying this formula

$$\begin{array}{ccccccc}
 f & \longrightarrow & \frac{f'}{y} & \longrightarrow \dots \longrightarrow & \frac{f^{(n-1)}}{y^{n-1}} & \longrightarrow & \frac{f^{(n)}}{y^n} \\
 \downarrow & & \downarrow & & \downarrow & & \\
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 \end{array}$$

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 \downarrow & & & & & & \\
 \frac{f}{y^4} & & & & & &
 \end{array}$$

Inductively applying this formula

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 f & \longrightarrow & \frac{f'}{y} & \longrightarrow \dots \longrightarrow & \frac{f^{(n-1)}}{y^{n-1}} & \longrightarrow & \frac{f^{(n)}}{y^n} \\
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 \frac{f}{y^4} & & \frac{f'}{y^5} & & \frac{f^{(n-1)}}{y^{n+3}} & &
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f	\longrightarrow	$\frac{f'}{y}$	$\longrightarrow \dots \longrightarrow$	$\frac{f^{(n-1)}}{y^{n-1}}$	\longrightarrow	$\frac{f^{(n)}}{y^n}$
\downarrow		\downarrow		\downarrow		
$\frac{f}{y^2}$	\longrightarrow	$\frac{f'}{y^3}$	$\longrightarrow \dots \longrightarrow$	$\frac{f^{(n-1)}}{y^{n+1}}$	\longrightarrow	$\frac{f^{(n)}}{y^{n+2}}$
\downarrow		\downarrow		\downarrow		
\vdots		\vdots		\vdots		\vdots
\downarrow		\downarrow		\downarrow		
$\frac{f}{y^{2m}}$	\longrightarrow	$\frac{f'}{y^{1+2m}}$	$\longrightarrow \dots \longrightarrow$	$\frac{f^{(n-1)}}{y^{n-1+2m}}$	\longrightarrow	$\frac{f^{(n)}}{y^{n+2m}}$
\downarrow		\downarrow		\downarrow		
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Then the **estimates of the lower derivatives** follow automatically

$$f^{(k)}(y) = \mathcal{O}(|y|^{n-1-\delta+(n-k)}), \quad \text{as } |y| \rightarrow \infty, \quad k \in \{0, \dots, n-1\}.$$

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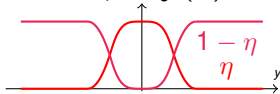
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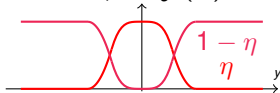


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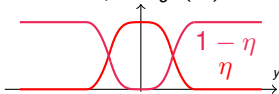
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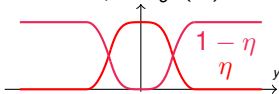
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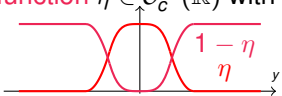
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J. Behrndt, P. S.: On a class of oscillatory integrals and their application to the time dependent Schrödinger equation. J. Math. Anal. and Appl. **543** (2025).

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- At $y = 0$: Splitting up the integral or multiplying cutoff

Application:

Superoscillations

The standard superoscillatory function

For $a > 1$ and $n \in \mathbb{N}$ consider

$$F_n(x) = \left(\cos\left(\frac{x}{n}\right) + ia \sin\left(\frac{x}{n}\right) \right)^n$$

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Moreover, for $\frac{|x|}{n} \ll 1$ these functions approximately behave as

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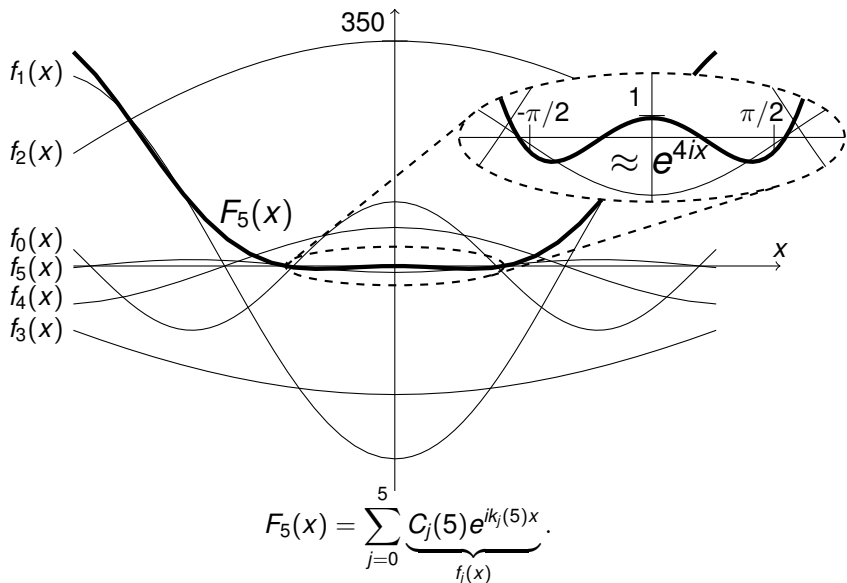
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In particular one can show that there converges

$$\lim_{n \rightarrow \infty} F_n(x) = e^{iax}.$$



Time evolution of superoscillations

Consider superoscillating initial conditions F_n of the time dependent Schrödinger equation

$$\begin{aligned} i \frac{\partial}{\partial t} \psi(t, x) &= \left(-\frac{\partial^2}{\partial x^2} + V(t, x) \right) \psi(t, x), & t > 0, x \in \mathbb{R}, \\ \psi(0, x) &= F_n(x), & x \in \mathbb{R}, \end{aligned}$$

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$$\begin{aligned} \lim_{n \rightarrow \infty} \Psi(t, x; F_n) &= \lim_{n \rightarrow \infty} \int_{\mathbb{R}} G(t, x, y) F_n(y) dy \\ &= \int_{\mathbb{R}} G(t, x, y) \lim_{n \rightarrow \infty} F_n(y) dy \\ &= \int_{\mathbb{R}} G(t, x, y) e^{ia y} dy = \Psi(t, x; e^{ia \cdot}). \end{aligned} \tag{2}$$

- The previous theory allows us to interchange the limit in (2).

Time evolution of superoscillations

Consider superoscillating initial conditions F_n of the time dependent Schrödinger equation

$$\begin{aligned}
 i \frac{\partial}{\partial t} \Psi(t, x) &= \left(-\frac{\partial^2}{\partial x^2} + V(t, x) \right) \Psi(t, x), & t > 0, x \in \mathbb{R}, \\
 \Psi(0, x) &= F_n(x), & x \in \mathbb{R},
 \end{aligned}$$

Is $\Psi_n(t, x; F_n)$ still superoscillating for $t > 0$?

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 &= \int_{\mathbb{R}} G(t, x, y) \lim_{n \rightarrow \infty} F_n(y) dy \\
 &= \int_{\mathbb{R}} G(t, x, y) e^{ia y} dy = \Psi(t, x; e^{ia \cdot}).
 \end{aligned} \tag{2}$$

- The previous theory allows us to interchange the limit in (2).
- Oscillatory behaviour of $\Psi(t, x; e^{ia \cdot})$ needs to be investigated.

Open questions

- Is the Green's function always of the form $G(t, x, y) = e^{iay^2} \tilde{G}(t, x, y)$?

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Thank you for your attention!